Complexity in Materials Science

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Outline

- **Introduction**

- **Pattern formation in physics: magnetic patterns as an example**

- **Multiscale structural complexity**

- **Self-induced glassiness and beyond: the role of frustration**

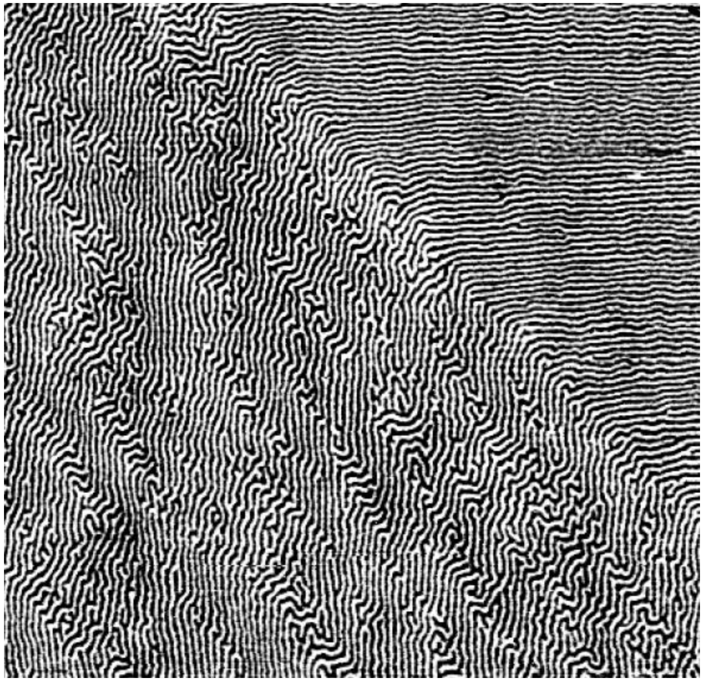
- **Experimental realization: elemental Nd**

- **Complexity of quantum frustrated systems**

- **Frustrations and complexity beyond materials science: machine learning, biological evolution and all that**

Intuitive feeling: crystals are simple, biological structures are complex

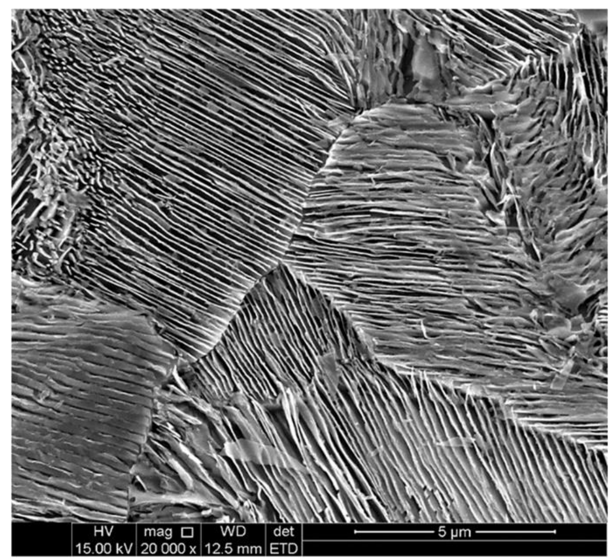
Complexity (“patterns”) in inorganic world



Stripe domains in ferromagnetic thin films

Stripes on a beach in tide zone

Microstructures in metals and alloys



Pearlitic structure in rail steel (Sci Rep 9, 7454 (2019))

Do we understand this? No, or, at least, not completely

Magnetic patterns

Example: strip domains in thin ferromagnetic films

PHYSICAL REVIEW B **69**, 064411 (2004)

Magnetization and domain structure of bcc $Fe_{81}Ni_{19}/Co$ (001) superlattices

R. Bručas, H. Hafermann, M. I. Katsnelson, I. L. Soroka, O. Eriksson, and B. Hjörvarsson

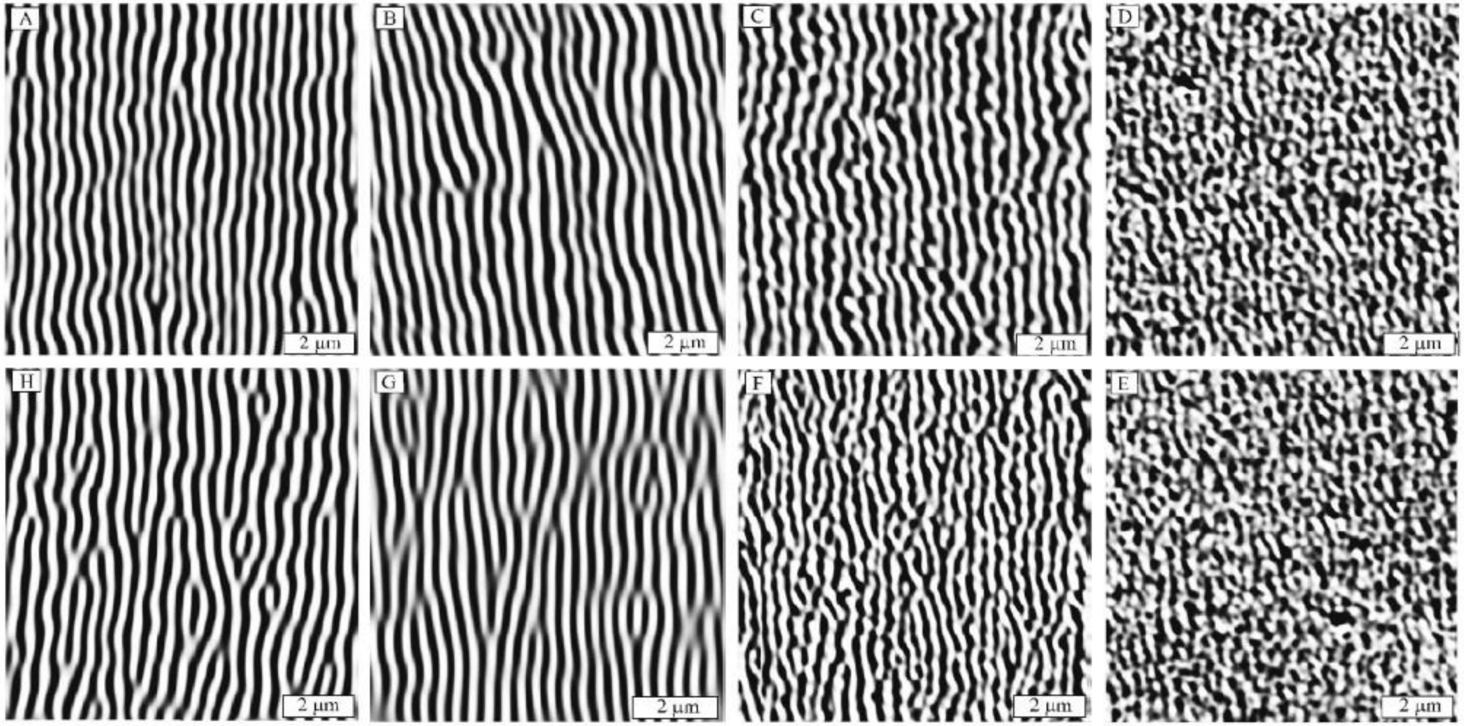


FIG. 2. The MFM images of the 420 nm thick $Fe_{81}Ni_{19}/Co$ superlattice at different externally applied in-plane magnetic fields: (a)—virgin (nonmagnetized) state; (b), (c), (d)—increasing field 8.3, 30, and 50 mT; (e), (f), (g)—decreasing field 50, 30, 8.3 mT; (h)—in remanent state.

Magnetic patterns III

Europhys. Lett., **73** (1), pp. 104–109 (2006)
DOI: 10.1209/epl/i2005-10367-8

Topological defects, pattern evolution, and hysteresis in thin magnetic films

P. A. Prudkovskii[1], A. N. Rubtsov[1] and M. I. Katsnelson[2]

$$H = \int \left(\frac{J_x}{2} \left(\frac{\partial \boldsymbol{m}}{\partial x} \right)^2 + \frac{J_y}{2} \left(\frac{\partial \boldsymbol{m}}{\partial y} \right)^2 - \frac{K}{2} m_z^2 - h m_y \right) \mathrm{d}^2 r + \\ + \frac{Q^2}{2} \int \int m_z(\boldsymbol{r}) \left(\frac{1}{|\boldsymbol{r} - \boldsymbol{r}'|} - \frac{1}{\sqrt{d^2 + (\boldsymbol{r} - \boldsymbol{r}')^2}} \right) m_z(\boldsymbol{r}') \mathrm{d}^2 r \mathrm{d}^2 r'.$$

Competition of exchange interactions (want homogeneous ferromagnetic state) and magnetic dipole-dipole interations (want total magnetization equal to zero)

Magnetic patterns IV

Classical Monte Carlo simulations

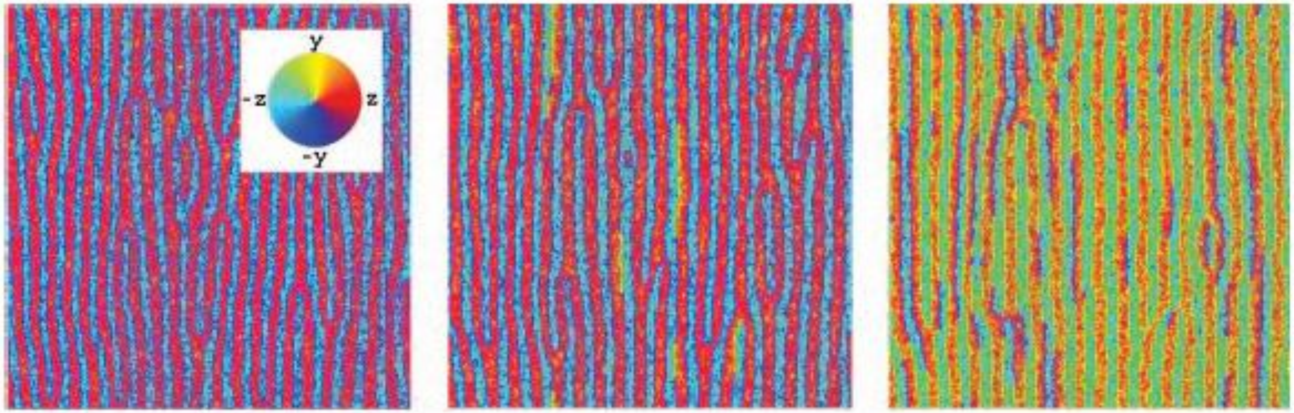


Fig. 2 – Snapshots of the stripe-domain system with the two-component order parameter at several points of the hysteresis loop for $\beta = 1$. The magnetic field is $h = 0$, $h = 0.3$, and $h = 0.6$, from left to right. The inset shows the color legend for the orientation of local magnetization.

We know the Hamiltonian and it is not very complicated

How to describe patterns and how to explain patterns?

What is complexity?

- Something that we immediately recognize when we see it, but very hard to define quantitatively

"I know it when I see it" (US Supreme Court Justice Potter Stewart, on obscenity)

- S. Lloyd, "Measures of complexity: a non-exhaustive list" – 40 different definitions
- Can be roughly divided into two categories:

- computational/descriptive complexities ("ultraviolet")
- effective/physical complexities ("infrared" or inter-scale)

Our definition: Multiscale structural complexity

Multi-scale structural complexity of natural patterns

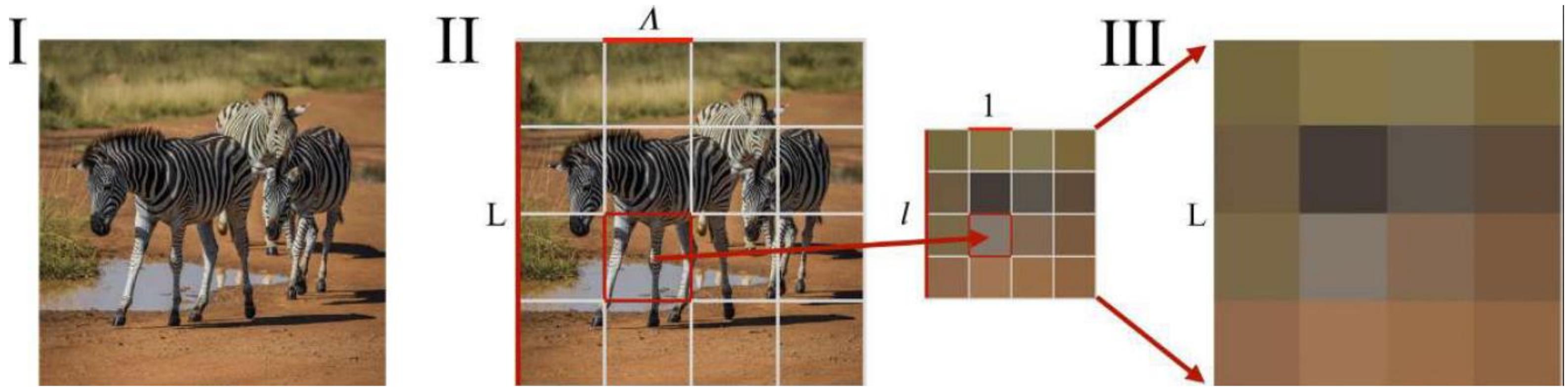
PNAS 117, 30241 (2020)

Andrey A. Bagrov[a,b,1,2], Ilia A. Iakovlev[b,1], Askar A. Iliasov[c], Mikhail I. Katsnelson[c,b], and Vladimir V. Mazurenko[b]

The idea: Complexity is dissimilarity at various scales

$f(x)$ multidimensional pattern $f_\Lambda(x)$ its coarse-grained version

Complexity is related to distances between $f_\Lambda(x)$ and $f_{\Lambda+d\Lambda}(x)$



$$\Delta_\Lambda = |\langle f_\Lambda(x)|f_{\Lambda+d\Lambda}(x)\rangle - \frac{1}{2}\left(\langle f_\Lambda(x)|f_\Lambda(x)\rangle + \langle f_{\Lambda+d\Lambda}(x)|f_{\Lambda+d\Lambda}(x)\rangle\right)| = \frac{1}{2}|\langle f_{\Lambda+d\Lambda}(x) - f_\Lambda(x)|f_{\Lambda+d\Lambda}(x) - f_\Lambda(x)\rangle|,$$

$$\langle f(x)|g(x)\rangle = \int_D dx\, f(x)g(x)$$

$$\mathcal{C} = \sum_\Lambda \frac{1}{d\Lambda}\Delta_\Lambda \to \int |\langle \frac{\partial f}{\partial\Lambda}|\frac{\partial f}{\partial\Lambda}\rangle| d\Lambda, \text{ as } d\Lambda \to 0$$

Multiscale structural complexity II

Solution of ink drop in water: Entropy should grow but complexity is not

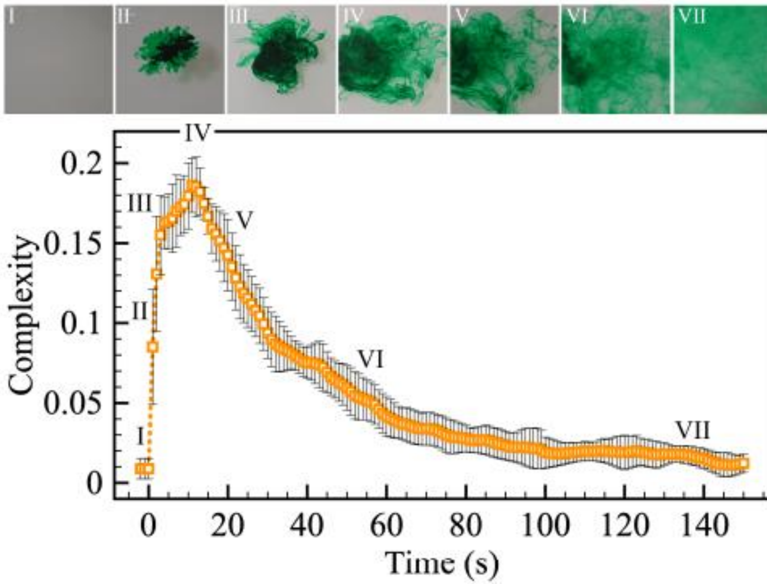


FIG. 7. The evolution of the complexity during the process of dissolving a food dye drop of 0.3 ml in water at 31°C.

Magnetic patterns

$$H = -J\sum_{nn'} \mathbf{S}_n\mathbf{S}_{n'} - \mathrm{D}\sum_{nn'}[\mathbf{S}_n \times \mathbf{S}_{n'}] - \sum_{n} BS_n^z$$

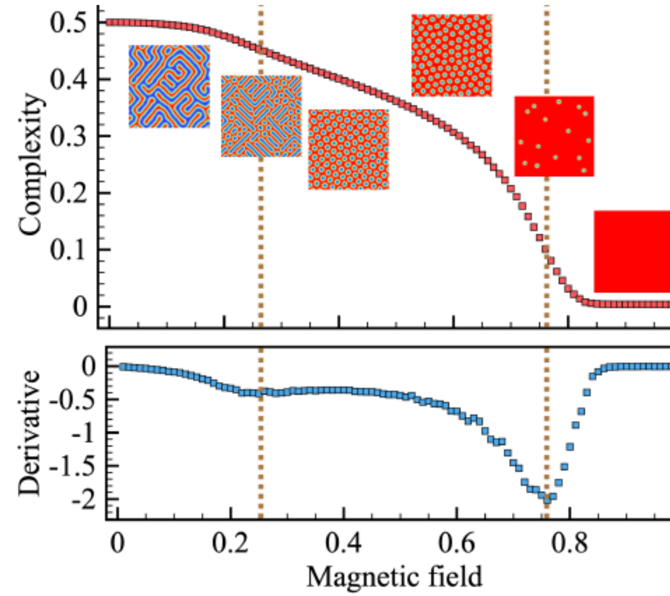


FIG. 4. (a) Magnetic field dependence of the complexity obtained from the simulations with spin Hamiltonian containing DM interaction with $J = 1$, $|\mathrm{D}| = 1$, $T = 0.02$. The error bars are smaller than the symbol size. (b) Complexity derivative we used for accurate detection of the phases boundaries.

And many other applications including biology and psychology

Nucleic Acids Research, 2024, **52**, 11045–11059
https://doi.org/10.1093/nar/gkae745
Advance access publication date: 28 August 2024
Genomics

Long range segmentation of prokaryotic genomes by gene age and functionality

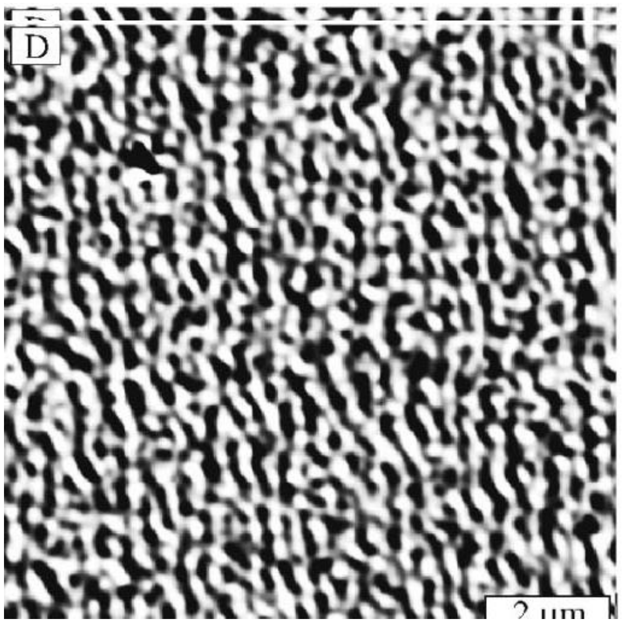
Yuri I. Wolf[1], Ilya V. Schurov[2], Kira S. Makarova[1], Mikhail I. Katsnelson[2] and Eugene V. Koonin[1,*]

Derivative detects changes of regime

Competing interactions and self-induced spin glasses

Special class of patterns: “chaotic” patterns

PHYSICAL REVIEW B **69**, 064411 (2004)



Hypothesis: a system wants to be modulated but cannot decide in which direction

$$E_m = \int\int d\mathbf{r}d\mathbf{r}' m(\mathbf{r})m(\mathbf{r}')\left[\frac{1}{|\mathbf{r}-\mathbf{r}'|} - \frac{1}{\sqrt{(\mathbf{r}-\mathbf{r}')^2+D^2}}\right]$$
$$= 2\pi \sum_{\mathbf{q}} m_{\mathbf{q}} m_{-\mathbf{q}} \frac{1-e^{-qD}}{q}, \tag{13}$$

where $m_{\mathbf{q}}$ is a two-dimensional Fourier component of the magnetization density. At the same time, the exchange energy can be written as

$$E_{exch} = \frac{1}{2}\alpha \sum_{\mathbf{q}} q^2 m_{\mathbf{q}} m_{-\mathbf{q}}, \tag{14}$$

so there is a finite value of the wave vector $q = q^*$ found from the condition

$$\frac{d}{dq}\left(2\pi\frac{1-e^{-qD}}{q} + \frac{1}{2}\alpha q^2\right) = 0 \tag{15}$$

Self-induced spin glasses II

PHYSICAL REVIEW B 93, 054410 (2016)

Stripe glasses in ferromagnetic thin films

Alessandro Principi* and Mikhail I. Katsnelson

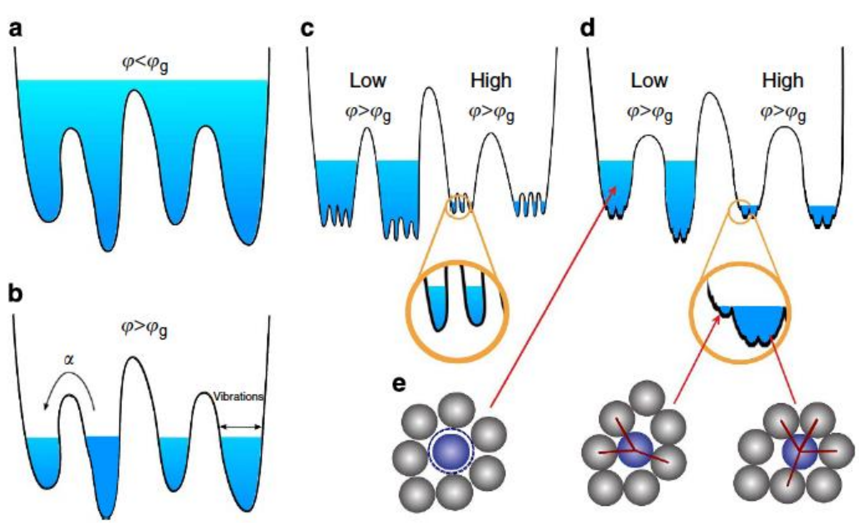
PRL 117, 137201 (2016) PHYSICAL REVIEW LETTERS week ending 23 SEPTEMBER 2016

Self-Induced Glassiness and Pattern Formation in Spin Systems Subject to Long-Range Interactions

Alessandro Principi* and Mikhail I. Katsnelson

Development of idea of stripe glass, J. Schmalian and P. G. Wolynes, PRL 2000

Glass: a system with an energy landscape characterizing by infinitely many local minima, with a broad distribution of barriers, relaxation at "any" time scale and aging (at thermal cycling you never go back to *exactly* the same state)



Picture from P. Charbonneau et al,

DOI: 10.1038/ncomms4725

Intermediate state between equilibrium and non-equilibrium, opportunity for history and memory ("stamp collection")

Self-induced spin glasses III

One of the ways to describe: R. Monasson, PRL 75, 2847 (1995)

$$\mathcal{H}_{\psi}[m,\lambda] = \mathcal{H}[m,\lambda] + g\int dr[m(r) - \psi(r)]^2$$

The second term describes attraction of our physical field $m(r)$ to some external field $\psi(r)$

If the system an be glued, with infinitely small interaction *g*, to macroscopically large number of configurations it should be considered as a glass

Then we calculate $F_g = \dfrac{\int \mathcal{D}\psi Z[\psi] F[\psi]}{\int \mathcal{D}\psi Z[\psi]}$ **and see whether the limits**

$F_{\text{eq}} = \lim_{N\to\infty} \lim_{g\to 0} F_g$ **and** $F = \lim_{g\to 0} \lim_{N\to\infty} F_g$ **are different**

If yes, this is self-induced glass

No disorder is needed (contrary to traditional view on spin glasses)

Self-induced spin glasses IV

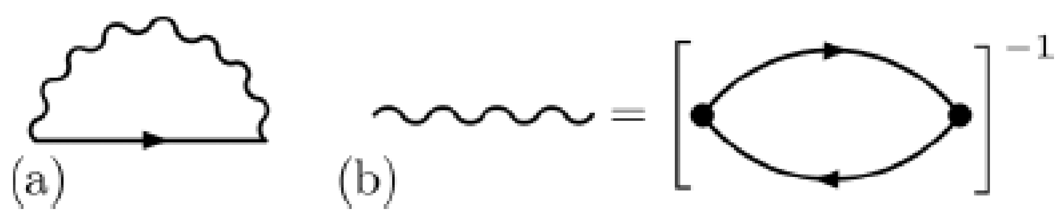
PHYSICAL REVIEW B **93**, 054410 (2016)

Stripe glasses in ferromagnetic thin films

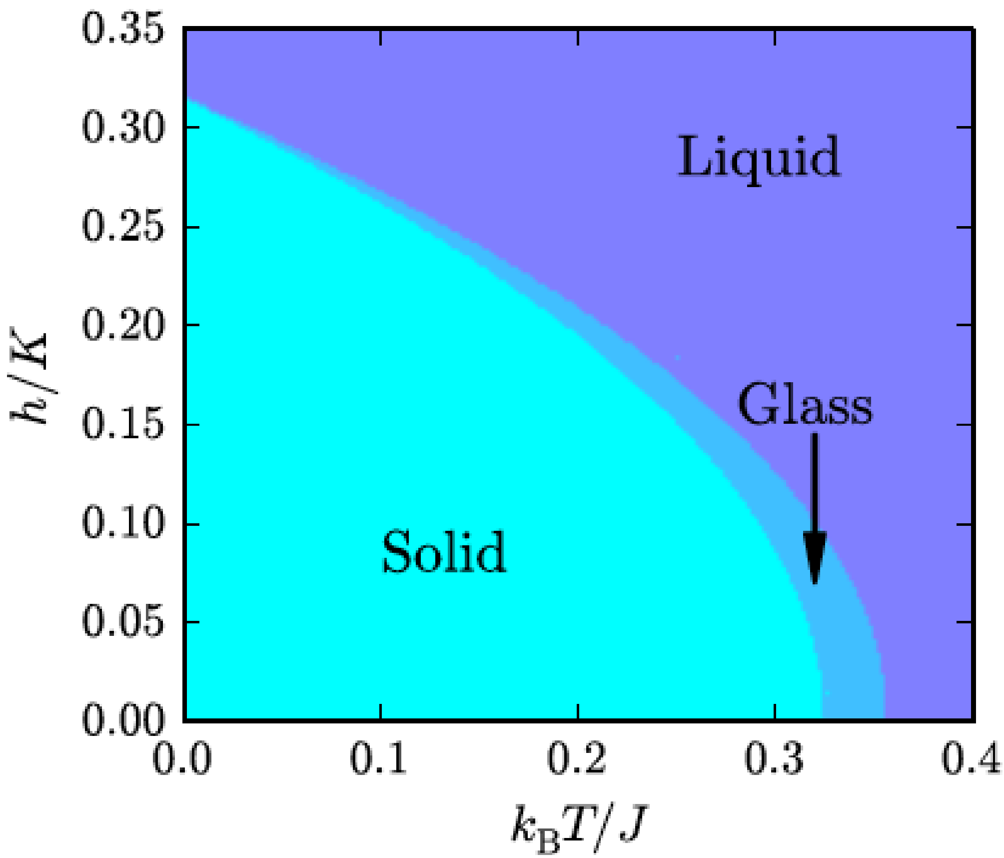
Alessandro Principi* and Mikhail I. Katsnelson

$$\mathcal{H}[\boldsymbol{m},\lambda] = \int d\boldsymbol{r}\{J[\partial_i m_j(\boldsymbol{r})]^2 - K m_z^2(\boldsymbol{r}) - 2\boldsymbol{h}(\boldsymbol{r})\cdot\boldsymbol{m}(\boldsymbol{r})\}$$
$$+ \frac{Q}{2\pi}\int d\boldsymbol{r}d\boldsymbol{r}' m_z(\boldsymbol{r})$$
$$\times\left[\frac{1}{|\boldsymbol{r}-\boldsymbol{r}'|} - \frac{1}{\sqrt{d^2+|\boldsymbol{r}-\boldsymbol{r}'|^2}}\right]m_z(\boldsymbol{r}')$$
$$+\int d\boldsymbol{r}\{\lambda(\boldsymbol{r})[\boldsymbol{m}^2(\boldsymbol{r})-1]\}. \quad (1)$$

Self-consistent screening approximation for spin propagators



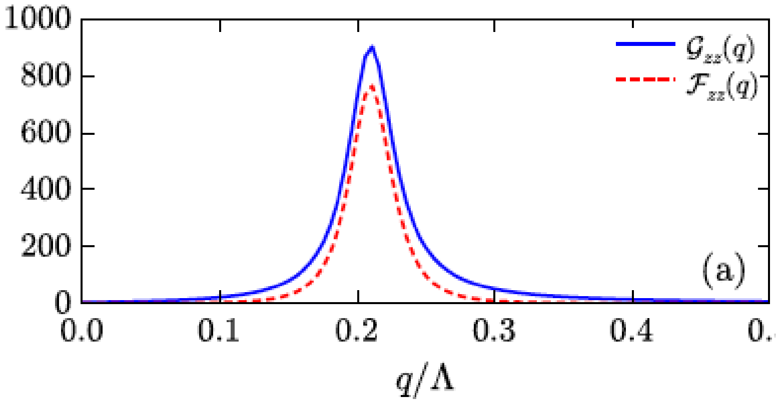
Self-induced spin glasses V



Phase diagram

Maximum at

$$q_0 \simeq [Q/(2J)]^{1/3} \neq 0$$



q-dependence of normal and anomalous ("glassy", non-ergodic spin-spin correlators

Glassiness without disorder?

Giorgio Parisi, Nobel Prize in physics 2021

"for the discovery of the interplay of disorder and fluctuations in physical systems from atomic to planetary scales."

Actually, disorder may be not needed, frustrations are enough (self-induced spin glass state in Nd)

Can we have something more or less exactly solvable?! – Yes!

PHYSICAL REVIEW B **109**, 144414 (2024)

Frustrated magnets in the limit of infinite dimensions: Dynamics and disorder-free glass transition

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(Received 16 November 2023; accepted 27 March 2024; published 18 April 2024)

The prototype theory: dynamical mean-field theory (DMFT) for strongly correlated systems (Metzner, Vollhardt, Georges, Kotliar and others)

Glassiness in infinite dimensions

Frustrations are necessary $$H = -\frac{1}{2}\sum_{i,j} J_{ij}^{\alpha\beta} S_i^{\alpha} S_j^{\beta} + \sum_i V(\mathbf{S}_i)$$

$$\mathbf{S}_i^2 = S_i^{\alpha} S_i^{\alpha} = 1$$

The limit of large dimensionality *d* $$J_{ij}^{\alpha\beta} = [f^{\alpha\beta}(\hat{t}/\sqrt{2d})]$$ e.g.

$$f^{\alpha\beta}(x) = J_0^{\alpha\beta} + J_1^{\alpha\beta} x + J_2^{\alpha\beta} x^2 + J_4^{\alpha\beta} x^4$$ means

$$J_{ij}^{\alpha\beta} = J_0^{\alpha\beta}\delta_{ij} + \frac{J_1^{\alpha\beta}}{\sqrt{2d}} t_{ij} + \frac{J_2^{\alpha\beta}}{2d}\sum_k t_{ik}t_{kj} + \frac{J_4^{\alpha\beta}}{4d^2}\sum_{k,l,m} t_{ik}t_{kl}t_{lm}t_{mj} \; .$$

The simplest frustrated model: $f^{\alpha\beta}(\varepsilon) = \delta^{\alpha\beta} f(\varepsilon)$ $\quad f(\varepsilon) = J(\varepsilon^2 - 1)$

Mean-field ordering temperature tends to zero at $d \to \infty$ in this model

Glassiness in infinite dimensions II

Cavity construction and mapping on effective single impurity

Purely dissipative Langevin dynamics

$$\dot{\mathbf{S}}_i = -\mathbf{S}_i \times (\mathbf{S}_i \times (\mathbf{N}_i + \boldsymbol{\nu}_i)) \\ = \mathbf{N}_i + \boldsymbol{\nu}_i - \mathbf{S}_i(\mathbf{S}_i \cdot (\mathbf{N}_i + \boldsymbol{\nu}_i))$$

$$\mathbf{N}_i = -\frac{\partial H}{\partial \mathbf{S}_i} = \mathbf{b}_i + \mathbf{F}_i \qquad b_i^\alpha = \sum_j J_{ij}^{\alpha\beta} S_j^\beta \qquad F^\alpha(\mathbf{S}_i) = -\partial V(\mathbf{S}_i)/\partial S_i^\alpha$$

$$\langle \nu_i^\alpha(t)\nu_j^\beta(t')\rangle = 2k_{\mathrm{B}}T\delta^{\alpha\beta}\delta_{ij}\delta(t-t')$$

Exactly mapped to a single-impurity dynamics with nonlocal in time "memory function"

Edwards-Anderson criterion of glassiness (local spin-spin correlation function tends to nonzero value in the limit of infinite time difference)

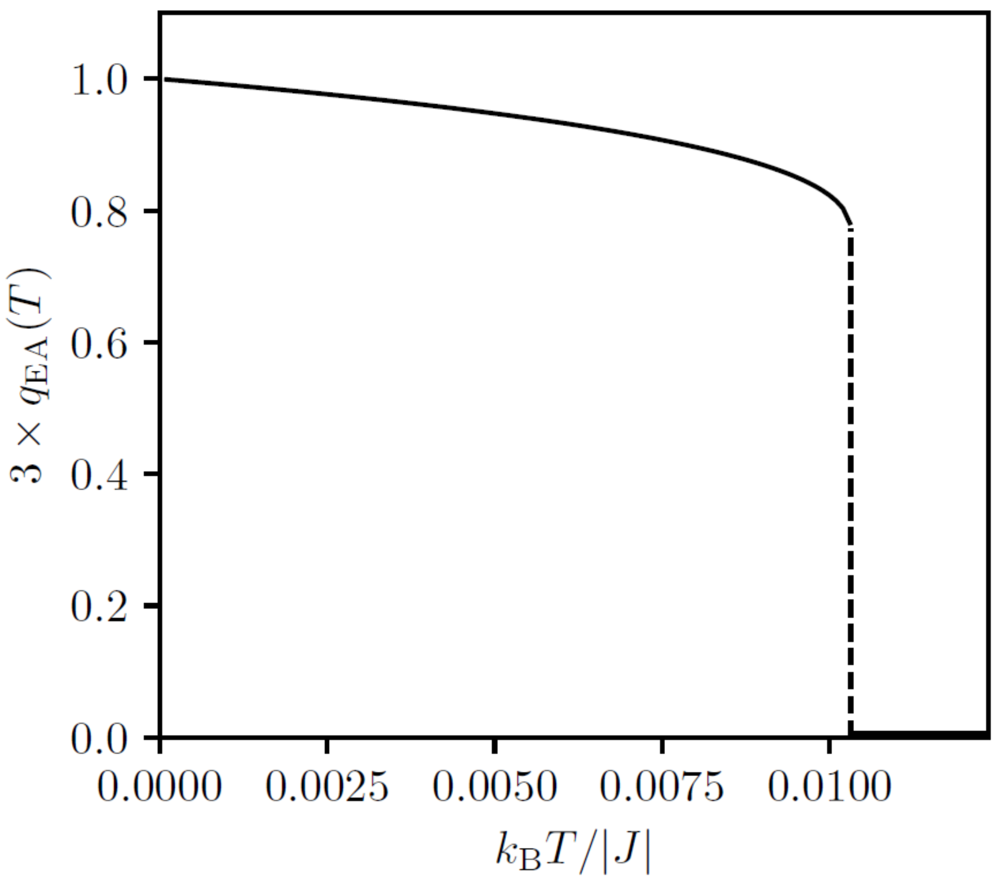
$$3q_{\mathrm{EA}}(T) = \lim_{|t-t'|\to\infty} \langle S^\alpha(t) S^\alpha(t')\rangle$$

Glassiness in infinite dimensions III

Isotropic model $f(\varepsilon) = J(\varepsilon^2 - 1)$

nonzero below the glass transition temperature $T_{\mathrm{g}} \simeq 0.0103|J|/k_{\mathrm{B}}$

First-order transition $q_{\mathrm{EA}}(T_{\mathrm{g}}) \simeq 0.2575$



Glassiness without disorder is theoretically possible!

Experimental observation of self-induced spin glass state: elemental Nd

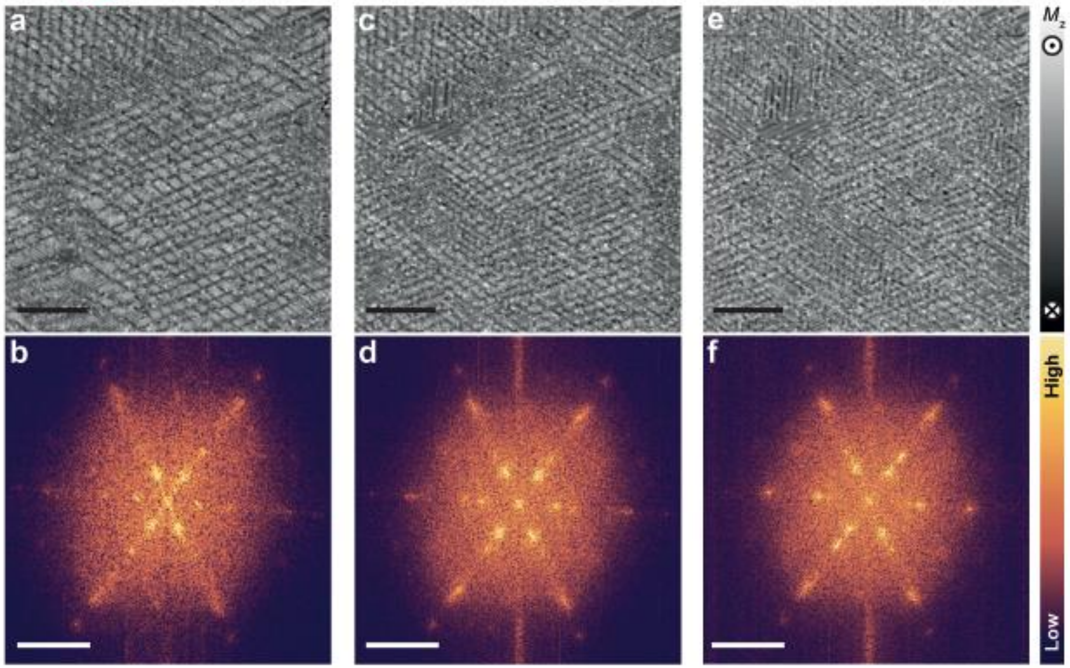
Self-induced spin glass state in elemental and crystalline neodymium

Umut Kamber, Anders Bergman, Andreas Eich, Diana Iuşan, Manuel Steinbrecher, Nadine Hauptmann, Lars Nordström, Mikhail I. Katsnelson, Daniel Wegner*, Olle Eriksson, Alexander A. Khajetoorians*

Science **368**, 966 (2020)

Spin-polarized STM experiment, Radboud University

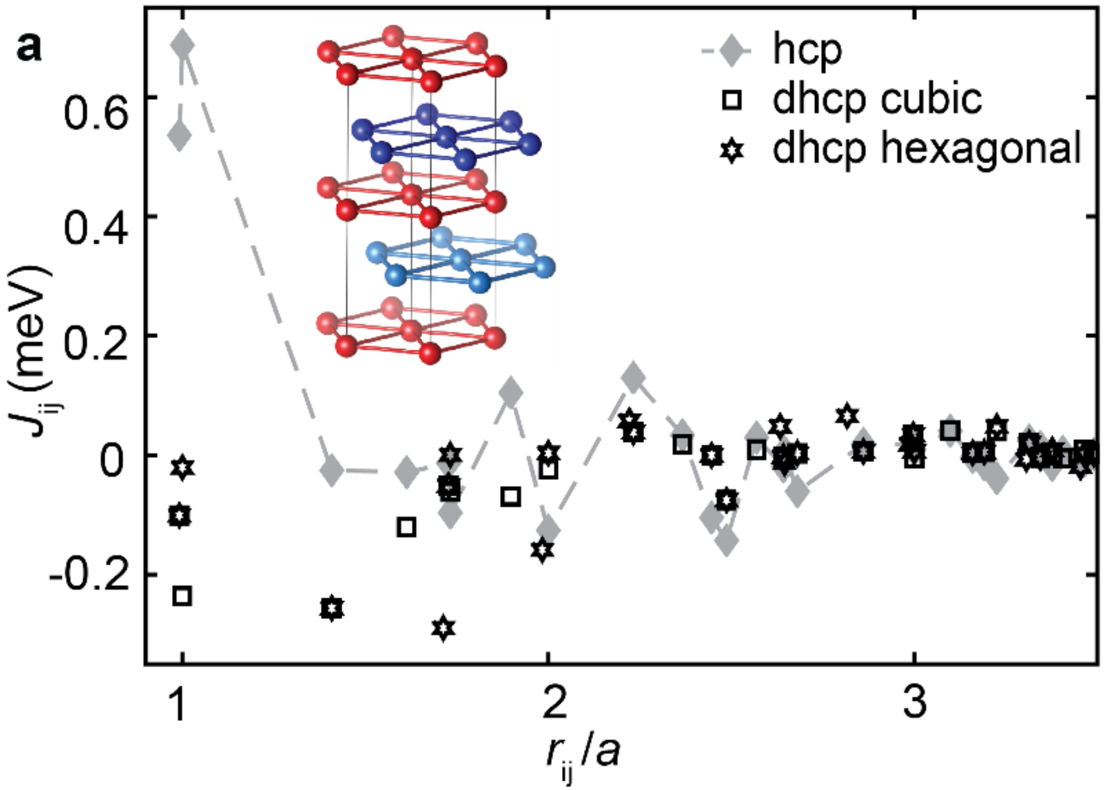
Magnetic structure: local correlations



The most important observation: aging. At thermocycling (or cyling magnetic field) the magnetic state is not exactly reproduced

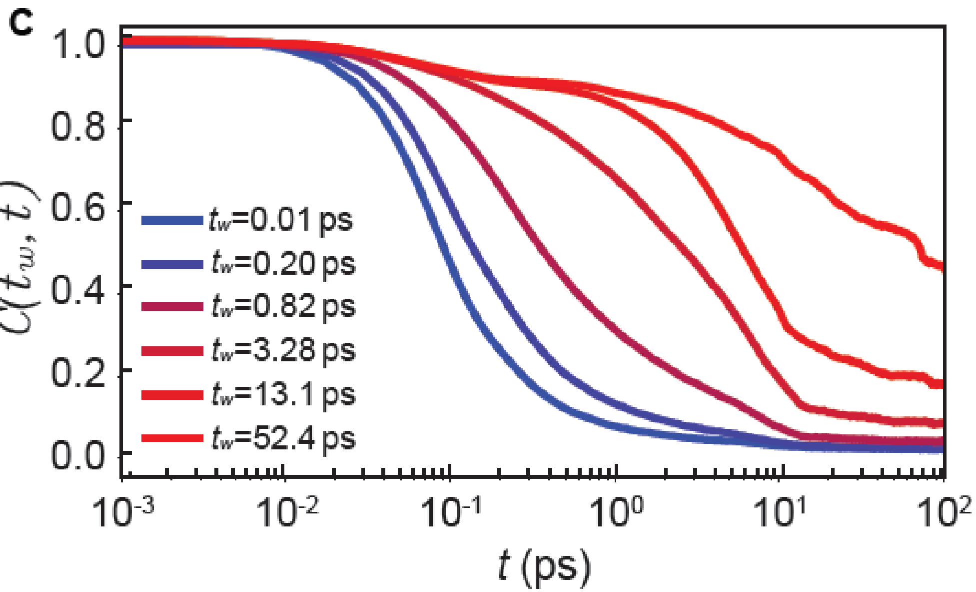
Ab initio: magnetic interactions in bulk Nd

Method: magnetic force theorem (Lichtenstein, Katsnelson, Antropov, Gubanov JMMM 1987)
Calculations: Uppsala team (Olle Eriksson group)



- Dhcp structure drives competing AFM interactions
- Frustrated magnetism

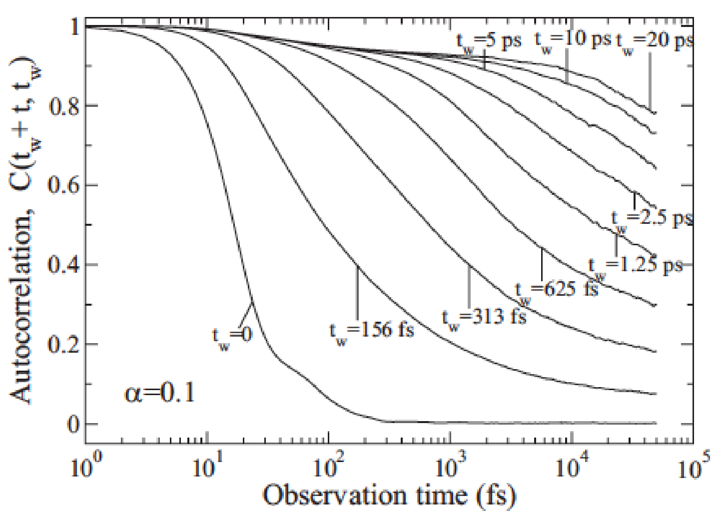
Spin-glass state in Nd: spin dynamics



Atomistic spin dynamics simulations

Typically spin-glass behavior

Autocorrelation function $C(t_w, t) = \langle \boldsymbol{m}_i(t + t_w) \cdot \boldsymbol{m}_i(t_w) \rangle$ for dhcp Nd at T = 1 K



To compare: the same for prototype *disordered* spin-glass Cu-Mn

B. Skubic et al, PRB 79, 024411 (2009)

Order from disorder

Thermally induced magnetic order from glassiness in elemental neodymium

NATURE PHYSICS | VOL 18 | AUGUST 2022 | 905–911

Benjamin Verlhac[1], Lorena Niggli[1], Anders Bergman[2], Umut Kamber[1], Andrey Bagrov[1,2], Diana Iuşan[2], Lars Nordström[2], Mikhail I. Katsnelson[1], Daniel Wegner[1], Olle Eriksson[2,3] and Alexander A. Khajetoorians[1]

Glassy state at low T and long-range order at T increase

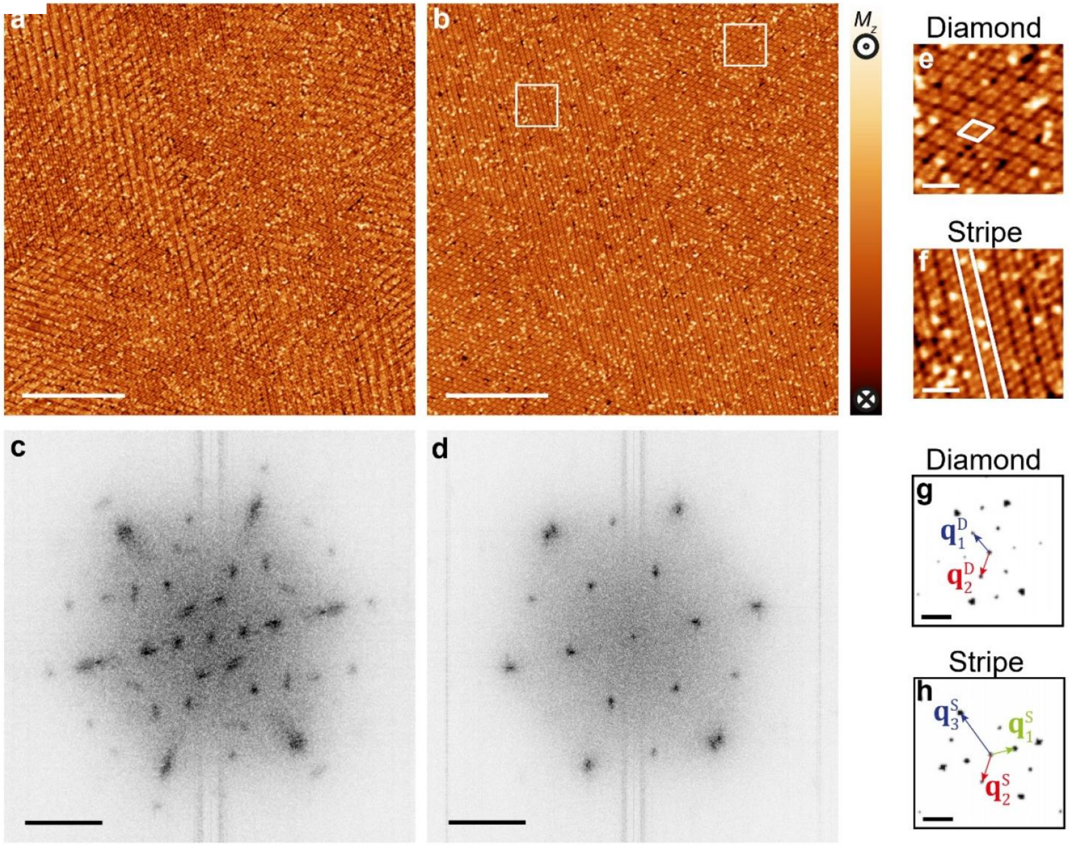
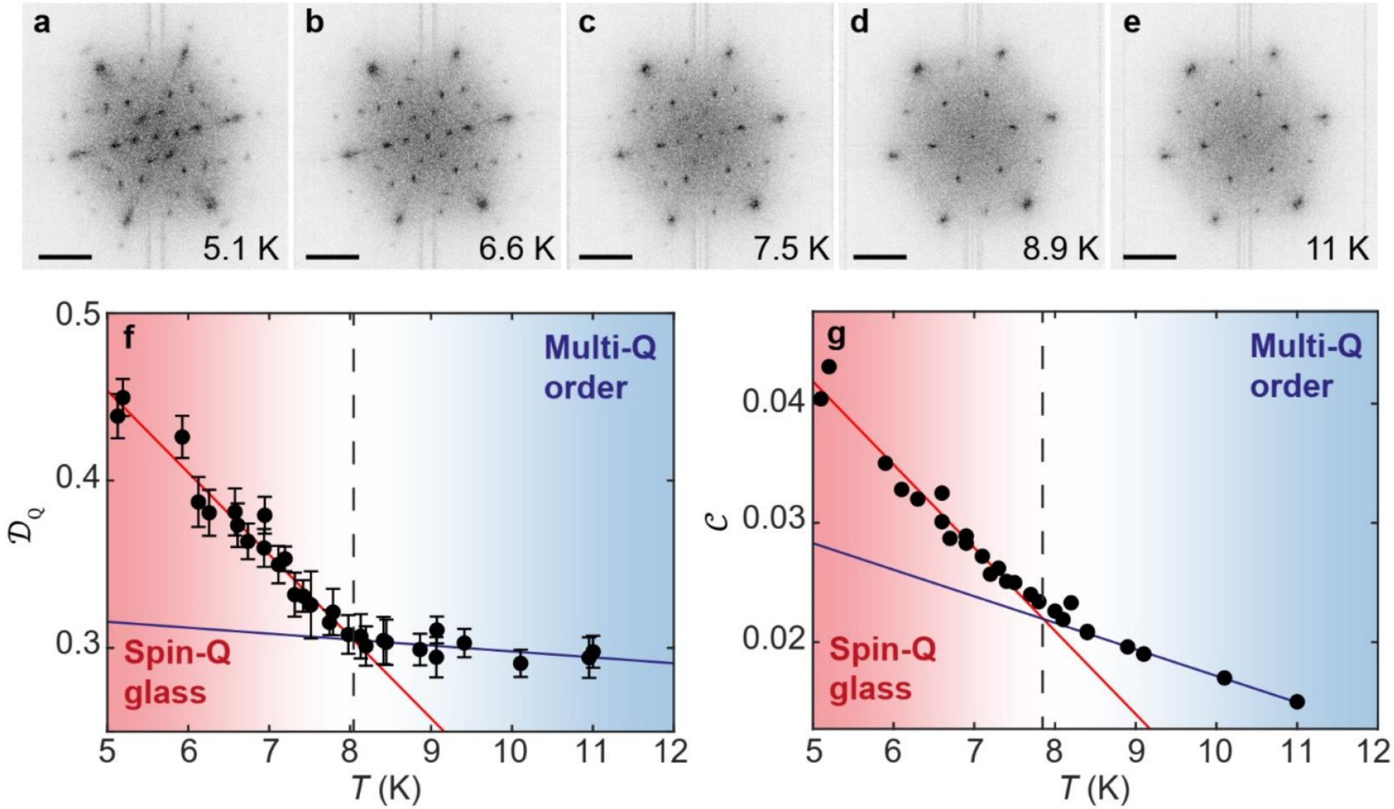


Figure 2: Emergence of long-range multi-Q order from the spin-Q glass state at elevated temperature. a,b. Magnetization images of the same region at T = 5.1 K and 11 K, respectively (I_t = 100 pA, a-b, scale bar: 50 nm). c,d. Corresponding Q-space images (scale bars: 3 nm^{-1}), illustrating the changes from strong local (i.e. lack of long-range) Q order toward multiple large-scale domains with well-defined long-range multi-Q order. e,f. Zoom-in images of the diamond-like (e) and stripe-like (f) patterns (scale bar: 5 nm). The locations of these images is shown by the white squares in b. g,h. Display of multi-Q state maps of the two apparent domains in the multi-Q ordered phase, where (g)

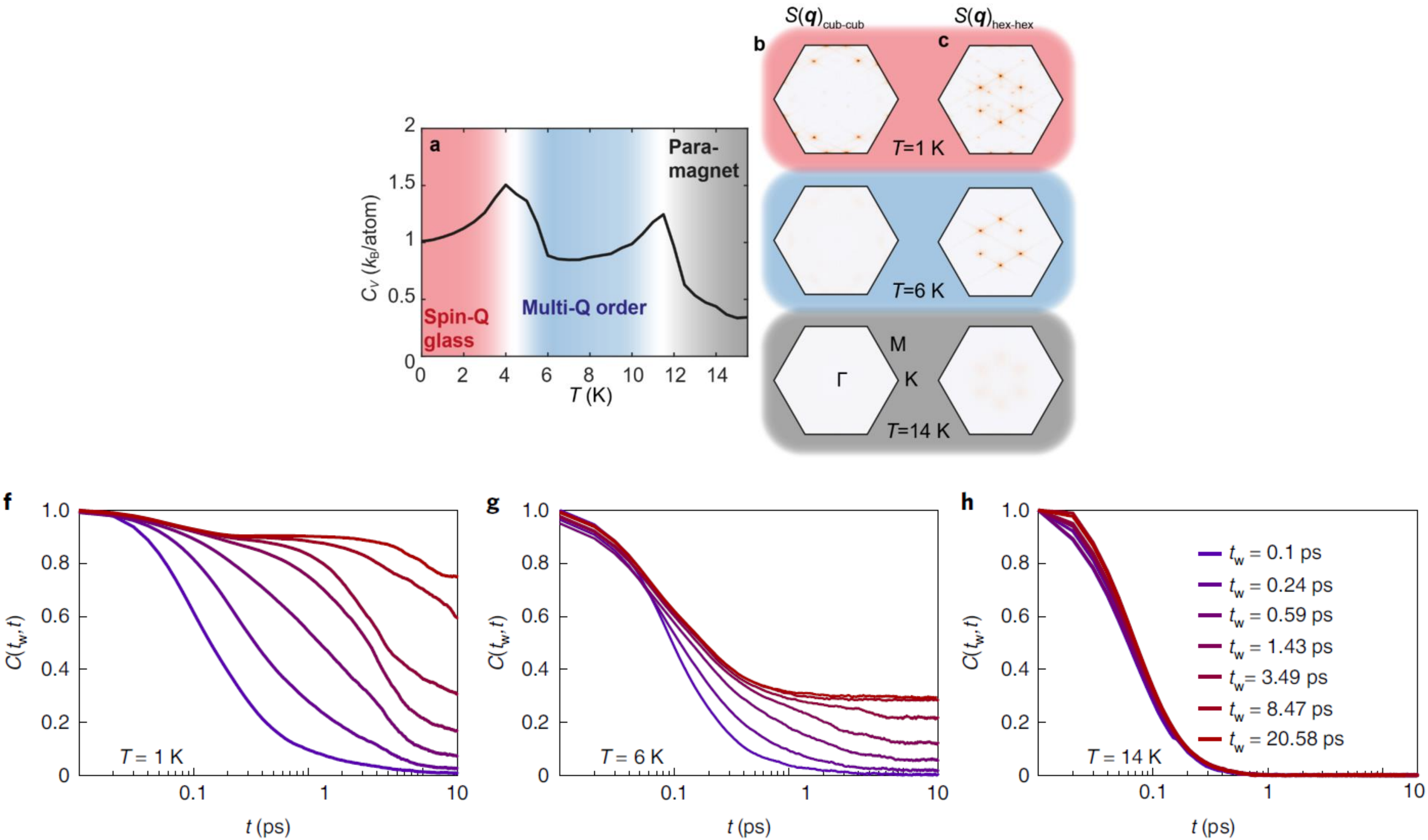
T=5K (a,c): spin glass
T=11K(b,d): (noncollinear) AFM

Order from disorder II



Phase transition at approx. 8K (seen via "complexity" measures) – right one is our multiscale structural complexity!

Order from disorder III



Theory: Atomistic simulations

Frustrations and complexity: Quantum case

Generalization properties of neural network approximations to frustrated magnet ground states

NATURE COMMUNICATIONS | (2020)11:1593

Tom Westerhout[1], Nikita Astrakhantsev[2,3,4], Konstantin S. Tikhonov[5,6,7], Mikhail I. Katsnelson[1,8] & Andrey A. Bagrov[1,8,9]

How to find true ground state of the quantum system?

In general, a very complicated problem (difficult to solve even for quantum computer!)

Idea: use of variational approach and train neural network to find "the best" trial function (G. Carleo and M. Troyer, Science 355, 602 (2017))

$$|\Psi_{GS}\rangle = \sum_{i=1}^{K} \psi_i |\mathcal{S}_i\rangle = \sum_{i=1}^{K} s_i |\psi_i| |\mathcal{S}_i\rangle$$

Generalization problem: to train NN for relatively small basis (*K* much smaller than total dim. of quantum space) and find good approximation to the true ground state

Frustrations and complexity: Quantum case II

Quantum S=1/2 Hamiltonian
NN and NNN interactions

$$\hat{H} = J_1 \sum_{\langle a,b \rangle} \hat{\boldsymbol{\sigma}}_a \otimes \hat{\boldsymbol{\sigma}}_b + J_2 \sum_{\langle\langle a,b \rangle\rangle} \hat{\boldsymbol{\sigma}}_a \otimes \hat{\boldsymbol{\sigma}}_b$$

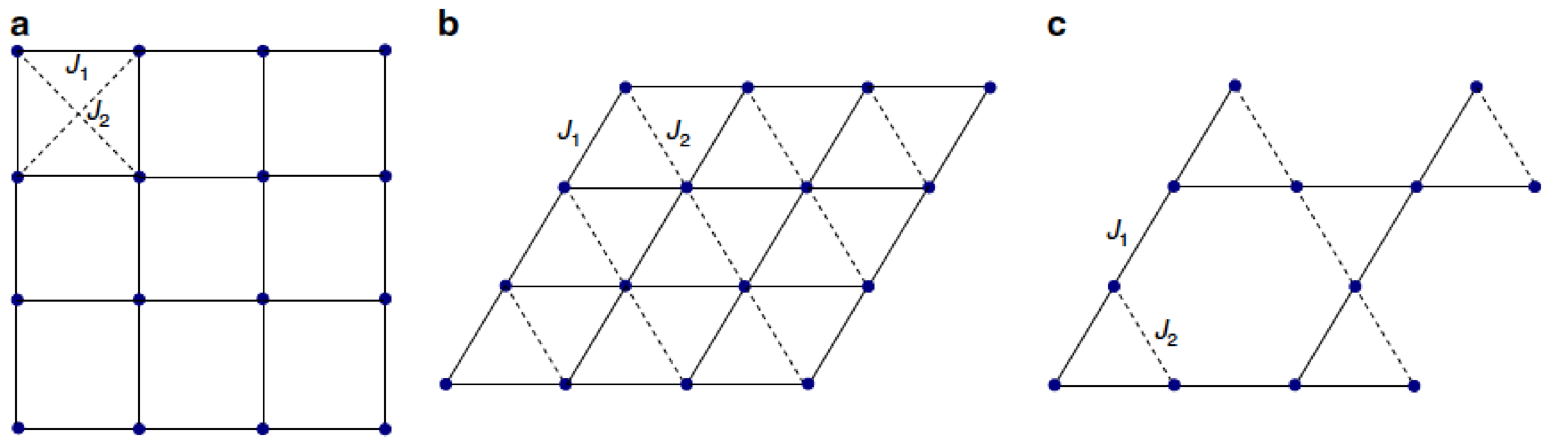


Fig. 1 Lattices considered in this work. We studied three frustrated antiferromagnetic Heisenberg models: **a** next-nearest neighbor J_1–J_2 model on square lattice; **b** anisotropic nearest-neighbor model on triangular lattice; **c** spatially anisotropic Kagome lattice. In all cases $J_2 = 0$ corresponds to the absence of frustration.

24 spins, dimensionality of Hilbert space $d = C_{12}^{24} \simeq 2.7 \cdot 10^6$

Still possible to calculate ground state exactly
Training for $K = 0.01\, d$ (small trial set)

Frustrations and complexity: Quantum case III

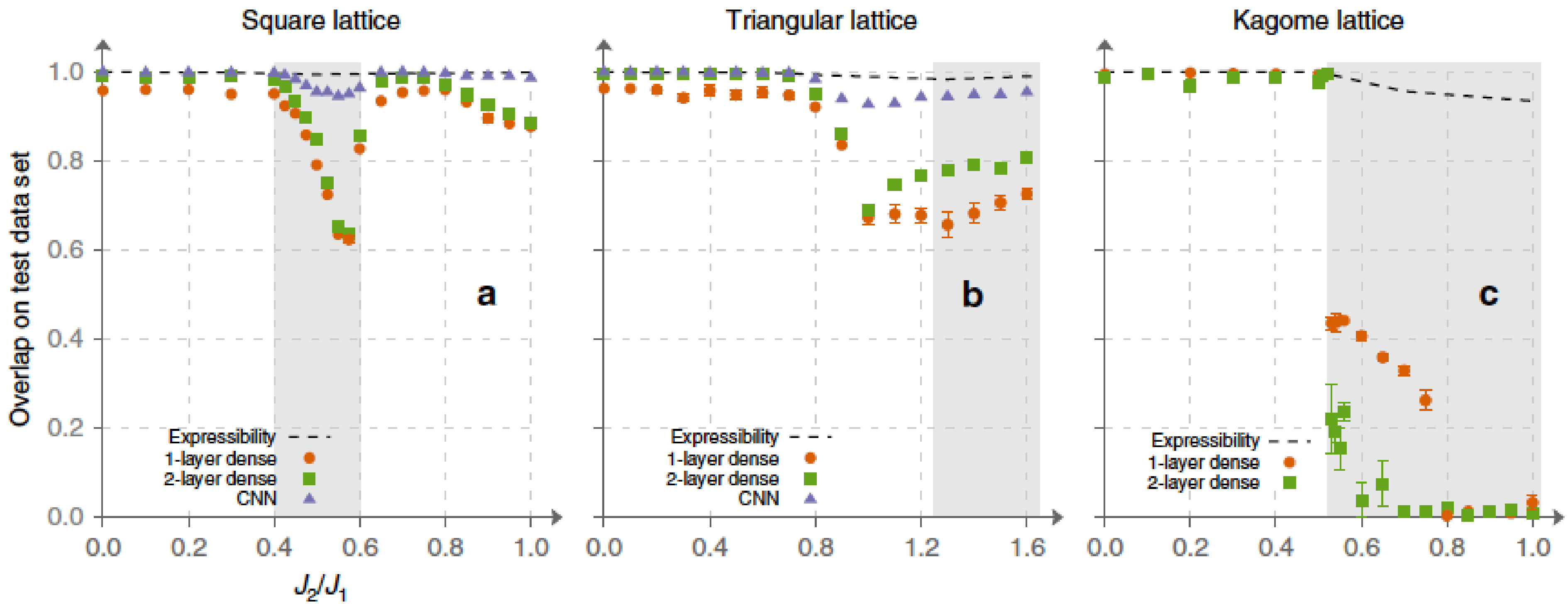


Fig. 2 Optimization results for 24-site clusters obtained with supervised learning and stochastic reconfiguration. Subfigures **a–c** were obtained using supervised learning of the sign structure. Overlap of the variational wave function with the exact ground state is shown as function of J_2/J_1 for square **a**, triangular **b**, and Kagome **c** lattices. Overlap was computed on the test dataset (not included into training and validation datasets). Note that generalization is poor in the frustrated regions (which are shaded on the plots). 1-layer dense, 2-layer dense, and convolutional neural network (CNN) architectures are described in Supplementary Note 1. Subfigures **d–f** show overlap between the variational wave function optimized using Stochastic Reconfiguration and the exact ground state for square, triangular, and Kagome lattices, respectively. Variational wave function was represented by two two-layer dense networks. A correlation between generalization quality and accuracy of the SR method is evident. On this figure, as well as on all the subsequent ones (both in the main text and Supplementary Notes 1 and 2), error bars represent standard error (SE) obtained by repeating simulations multiple times.

Frustrations and complexity: Quantum case IV

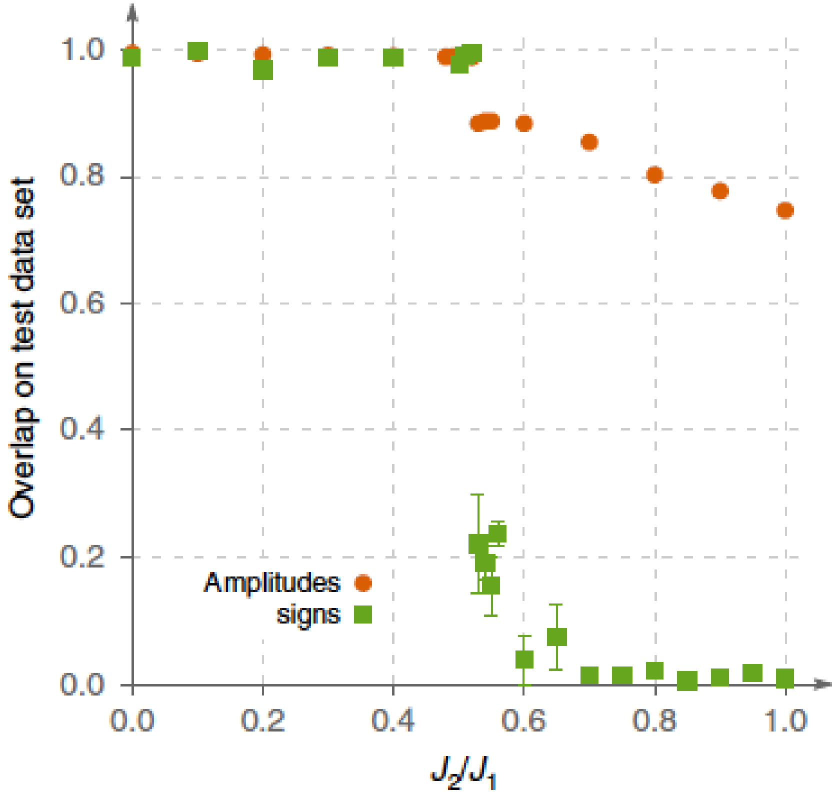


Fig. 4 Generalization of signs and amplitudes. We compare generalization quality as measured by overlap for learning the sign structure (red circles) and amplitude structure (green squares) for 24-site Kagome lattice for two-layer dense architecture. Note that both curves decrease in the frustrated region, but the sign structure is much harder to learn.

It is *sign* structure which is difficult to learn in frustrated case!!!

Relation to sign problem in QMC?!

"Somehow it seems to fill my head with ideas –only I don't exactly know what they are!" (Through the Looking-Glass, and What Alice Found There)

Further development

Many-body quantum sign structures as non-glassy Ising models

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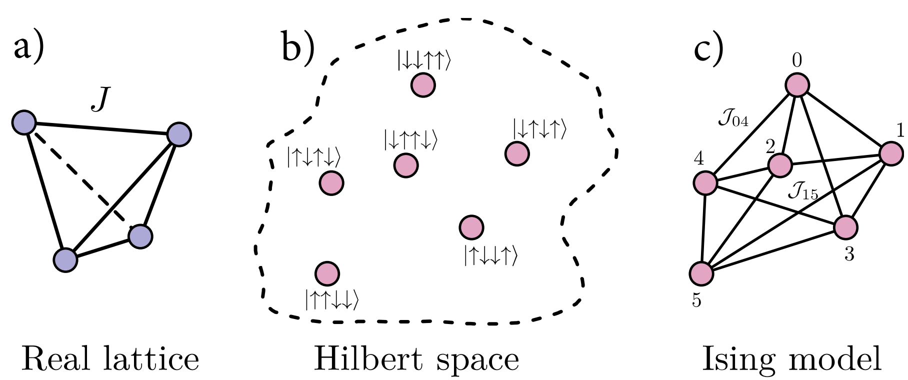
The idea: use machine learning to find amplitudes and then map onto efficient Ising model

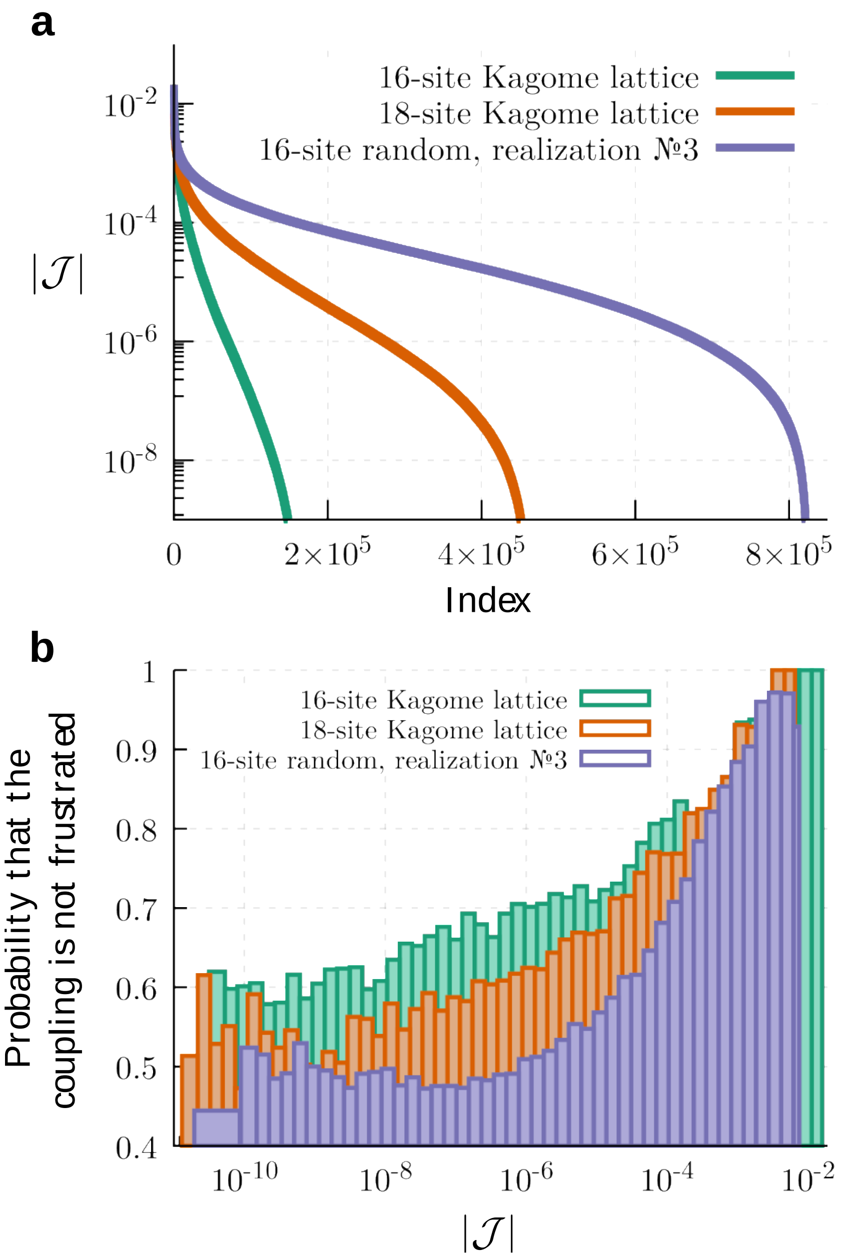
$$|\Psi_{\mathrm{GS}}\rangle = \sum_{i=1}^{K} \psi_i |\mathcal{S}_i\rangle = \sum_{i=1}^{K} s_i |\psi_i| |\mathcal{S}_i\rangle$$

When amplitudes are known the trial ground state energy $<\Psi|H|\Psi>$

is a bilinear function of signs s_i, and we have Ising optimization problem in K-dimensional space; K is very big but it turns out that the model is not glassy and can be optimized without too serious problems

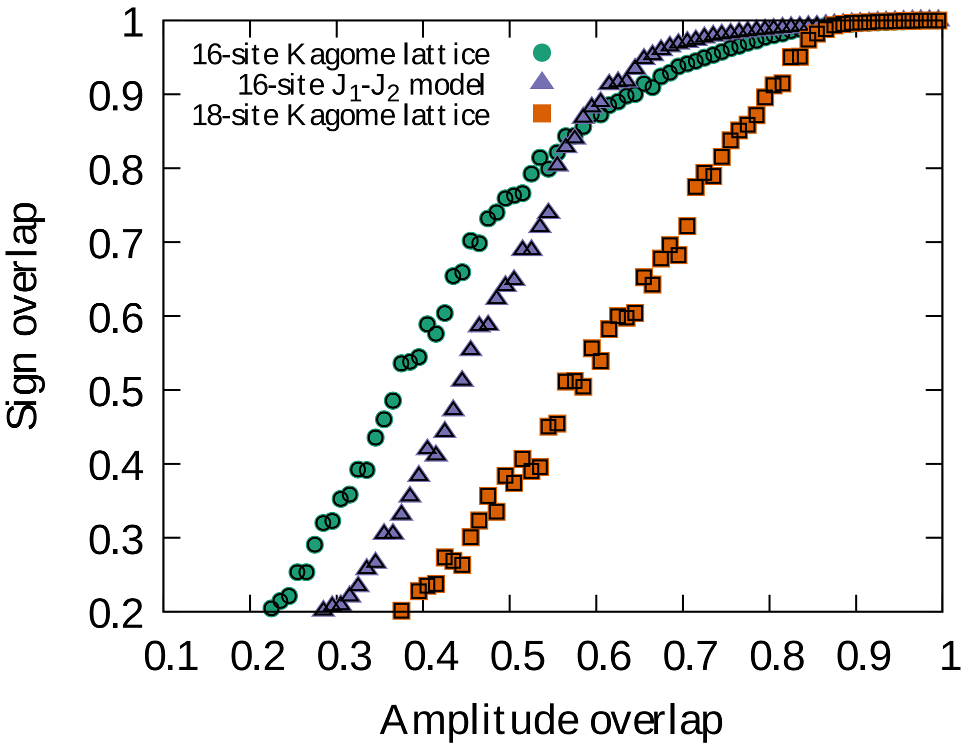


Further development II



It turns out that even for initially frustrated quantum spin models the effective Ising model is not frustrated, both couplings are small and optimization is quite efficient

Further development III



The quality of optimization is quite robust with respect to uncertainties in amplitudes (overlap with the exact ground state)

Analogies with biological evolution?

Toward a theory of evolution as multilevel learning

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Thermodynamics of evolution and the origin of life

Vitaly Vanchurin[a,b,1], Yuri I. Wolf[a], Eugene V. Koonin[a,1], and Mikhail I. Katsnelson[c,1]

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Table 1. Corresponding quantities in thermodynamics, machine learning, and evolutionary biology

	Thermodynamics	Machine learning	Evolutionary biology
$\mathbf{x}$	Microscopic physical degrees of freedom	Variables describing training dataset (nontrainable variables)	Variables describing environment
$\mathbf{q}$	Generalized coordinates (e.g., volume)	Weight matrix and bias vector (trainable variables)	Trainable variables (genotype, phenotype)
$H(\mathbf{x},\mathbf{q})$	Energy	Loss function	Additive fitness, $H(x,q) = -T\log f(q)$
$S(\mathbf{q})$	Entropy of physical system	Entropy of nontrainable variables	Entropy of biological system
$U(\mathbf{q})$	Internal energy	Average loss function	Average additive fitness
$Z(T,\mathbf{q})$	Partition function	Partition function	Macroscopic fitness
$F(T,\mathbf{q})$	Helmholtz free energy	Free energy	Adaptive potential (macroscopic additive fitness)
$\Omega(T,\mu)$	Grand potential, $\Omega_p(\mathscr{T},\mathscr{M})$	Grand potential	Grand potential, $\Omega_b(T,\mu)$
T or $\mathscr{T}$	Physical temperature, $\mathscr{T}$	Temperature	Evolutionary temperature, T
μ or $\mathscr{M}$	Chemical potential, $\mathscr{M}$	Absent in conventional machine learning	Evolutionary potential, μ
N_e or N	Number of molecules, N	Number of neurons, N	Effective population size, N_e
K	Absent in conventional physics	Number of trainable variables	Number of adaptable variables

Energy landscape in physics is similar to fitness landscape in biology

Analogies with biological evolution II

Can the change of e.g. biological temperature switch fitness landscape from a few well-defined peaks to a glassy-like with many directions of possible evolution?

Explaining the Cambrian "Explosion" of Animals

Charles R. Marshall

Annu. Rev. Earth Planet. Sci. 2006. 34:355–84

Australian Journal of Zoology
http://dx.doi.org/10.1071/ZO13052

The evolution of morphogenetic fitness landscapes: conceptualising the interplay between the developmental and ecological drivers of morphological innovation

Charles R. Marshall

Cambrian Exposion as an analog of magnetic phase transitions in neodymium?!

Well… for me (as a physicist) it is a good place to stop

What remains beyond the talk (examples)?

Complexity in metallurgy, interplay of magnetism and structural state in steel

PHYSICAL REVIEW B **90**, 094101 (2014)

Role of magnetic degrees of freedom in a scenario of phase transformations in steel

I. K. Razumov,[1,2,*] D. V. Boukhvalov,[3] M. V. Petrik,[2] V. N. Urtsev,[4] A. V. Shmakov,[4] M. I. Katsnelson,[5,6] and Yu. N. Gornostyrev[1,2]

PHYSICAL REVIEW APPLIED **7**, 014002 (2017)

Autocatalytic Mechanism of Pearlite Transformation in Steel

I. K. Razumov,[1,2,*] Yu. N. Gornostyrev,[1,2,4] and M. I. Katsnelson[3,4]

Perspectives to use complex patterns for neuromorphic computations etc. – e.g. adatoms at black P

ARTICLE
DOI: 10.1038/s41467-018-06337-4 OPEN

An orbitally derived single-atom magnetic memory

Brian Kiraly[1], Alexander N. Rudenko[1,2,3], Werner M.J. van Weerdenburg[1], Daniel Wegner[1], Mikhail I. Katsnelson[1] & Alexander A. Khajetoorians[1]

PHYSICAL REVIEW LETTERS **128**, 106801 (2022)

Gating Orbital Memory with an Atomic Donor

Elze J. Knol, Brian Kiraly, Alexander N. Rudenko, Werner M. J. van Weerdenburg, Mikhail I. Katsnelson, and Alexander A. Khajetoorians

LETTERS
https://doi.org/10.1038/s41565-020-00838-4
nature nanotechnology
Check for updates

An atomic Boltzmann machine capable of self-adaption

Brian Kiraly[1,3], Elze J. Knol[1,3], Werner M. J. van Weerdenburg[1], Hilbert J. Kappen[2] and Alexander A. Khajetoorians[1]✉

MANY THANKS FOR YOUR ATTENTION