

# Structure–Composition–Property Relationships in Antiperovskite Nitrides: Guiding a Rational Alloy Design

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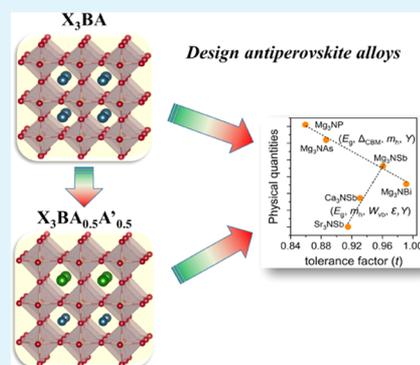
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**ABSTRACT:** The alloy strategy through the A- or X-site is a common method for experimental preparation of high-performance and stable lead-based perovskite solar cells. As one of the important candidates for lead-free and stable photovoltaic absorbers, the inorganic antiperovskite family has recently been reported to exhibit excellent optoelectronic properties. However, the current reports on the design of antiperovskite alloys are rare. In this work, we investigated the previously overlooked electronic property (e.g., conduction band convergence), static dielectric constant, and exciton binding energy in inorganic antiperovskite nitrides by first-principles calculations. Then, we revealed a linear relationship between the tolerance factor and various physical quantities. Guided by the established structure–composition–property relationship in six antiperovskite nitrides  $X_3NA$  ( $X^{2+} = Mg^{2+}, Ca^{2+}, Sr^{2+}$ ;  $A^{3-} = P^{3-}, As^{3-}, Sb^{3-}, Bi^{3-}$ ), for the first time, we designed a promising antiperovskite alloy  $Mg_3NAS_{0.5}Bi_{0.5}$  with a quasi-direct band gap of 1.402 eV. Finally, we made a comprehensive comparison between antiperovskite nitrides and conventional halide perovskites for pointing out the future direction for device applications.

**KEYWORDS:** antiperovskite, electronic property, structure–composition–property relationship, alloy design, first-principles calculations



## INTRODUCTION

Perovskite semiconductors have attracted intensive research interest in recent years due to their versatile compositions and tunable electronic and optical properties.<sup>1–3</sup> As an important member of the perovskite family, lead (Pb) halide perovskites  $APbX_3$  ( $A^+ = CH_3NH_3^+, CH(NH_2)_2^+, Cs^+$ ;  $X^- = I^-, Br^-, Cl^-$ ) have achieved great success in the application of solar cells,<sup>1</sup> light-emitting diodes,<sup>2</sup> and photodetectors<sup>3,4</sup> because of their intriguing optoelectronic properties (see Figure 1). For example, the power conversion efficiency (PCE) of perovskite solar cells has been rapidly improved from 3.8 to 25.5% within only one decade.<sup>5,6</sup> Despite the exciting progress, long-term instability and Pb toxicity are still key challenges hindering perovskite solar cells from practical applications.

An effective strategy to solve the above problem is to carry out a cation or anion substitution. For example, the most straightforward way to eliminate the toxic Pb(II) is to replace it with Sn(II). The obtained perovskites  $ASnX_3$  have a crystal structure and electronic properties similar to Pb counterparts, but unfortunately they suffer from serious instability issues (i.e., against the oxidation to  $Sn^{4+}$ ).<sup>8</sup> When replacing  $I^-$  with  $S^{2-}/Se^{2-}$ , it can lead to chalcogenide perovskites  $ABX_3$  ( $A^{2+} = Ca^{2+}, Sr^{2+}, Ba^{2+}$ ;  $B^{4+} = Ti^{4+}, Zr^{4+}, Hf^{4+}$ ;  $X^{2-} = S^{2-}, Se^{2-}$ ) with robust thermal stability and nontoxic elements (see Figure 1).<sup>9,10</sup> Due to the large electronegativity difference between B- and X-site ions and the disordered ground-state structure, they usually exhibit a large band gap, and they are difficult to form

high-quality films at low temperatures.<sup>9</sup> Although the experimentally synthesized  $Ba_3Zr_2S_7$  possesses a surprisingly optimal band gap of 1.28 eV, its valence band maximum (VBM) and conduction band minimum (CBM) do not exhibit sufficient dispersion (i.e., small band width  $W$  and large carrier effective mass  $m$ ).<sup>11</sup> Therefore, it is very challenging to realize well-performed solar cells based on chalcogenide perovskites.

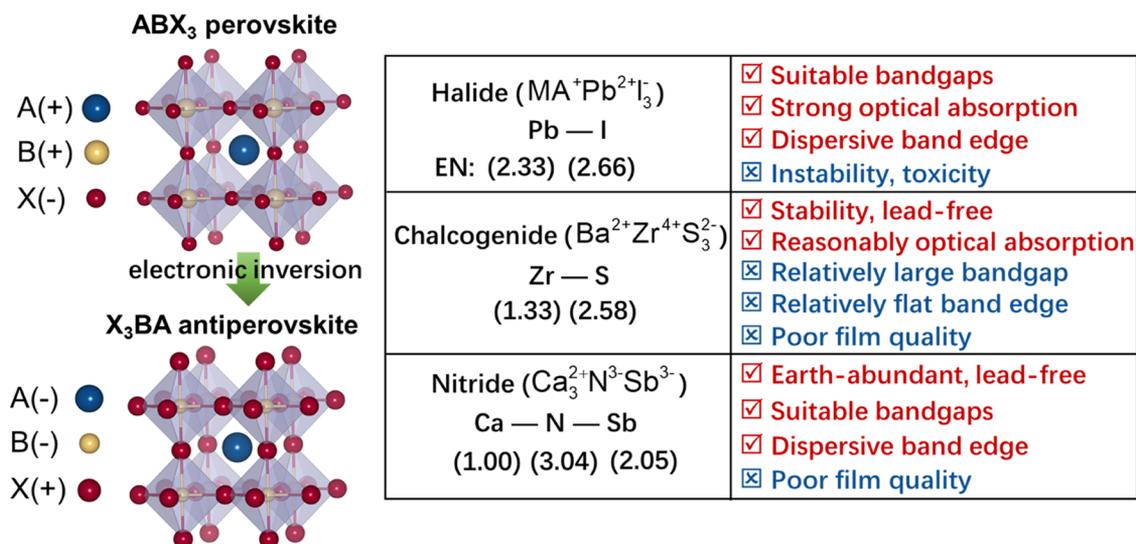
Another effective strategy is to reverse the ion type of perovskite lattice sites. Specifically, through electronically inverting the formula  $ABX_3$  (A and B are cations, and X is an anion), a class of antiperovskite  $X_3BA$  can be obtained, in which anions occupy the A and B sites and cations occupy the X sites (see Figure 1).<sup>12</sup> Recently, Gebhardt et al. proposed a series of inverse-hybrid perovskites  $(CH_3NH_3)_3BA$  (B: monovalent anions, A: divalent anions or B: divalent anions, A: monovalent anions) for photovoltaic applications. Unfortunately, these predicted compounds still do not exhibit the dispersive valence and conduction bands near the Fermi level.<sup>13,14</sup> At present, several synthesized hybrid organic–inorganic antiperovskites have shown potential applications in

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**Figure 1.** Crystal structures of cubic perovskites ABX<sub>3</sub> and antiperovskites X<sub>3</sub>BA. Three representatives, lead halide perovskite, chalcogenide perovskite, and antiperovskite nitride, are presented, with their corresponding preponderances and disadvantages. The electronegativity (EN) values of the elements are given in parentheses (Pauling Units).

the field of molecular ferroelectric,<sup>15</sup> but no hybrid antiperovskite suitable as a photovoltaic absorber has been experimentally reported. On the contrary, all-inorganic antiperovskite semiconductors (i.e., X<sub>3</sub>NA, X<sup>2+</sup> = Mg<sup>2+</sup>, Ca<sup>2+</sup>, Sr<sup>2+</sup>, Ba<sup>2+</sup>; A<sup>3-</sup> = P<sup>3-</sup>, As<sup>3-</sup>, Sb<sup>3-</sup>, Bi<sup>3-</sup>) have been reported to exhibit attractive electronic structures, optical absorptions, ion migration barriers, and defect properties, which show great potential applications in photovoltaics, thermoelectrics, and solid electrolytes (see Figure 1).<sup>16–18</sup> For example, Dai<sup>17</sup> and Mochizuki<sup>18</sup> have successively proposed Ca<sub>3</sub>NSb, Ca<sub>3</sub>Nb, Mg<sub>3</sub>NP, and Sr<sub>3</sub>NP as promising photovoltaic absorbers because of their suitable direct band gaps and high optical absorption coefficients. Although the antiperovskite family greatly expands the composition space for searching Pb-free and stable photovoltaic candidates, there are few reports that establish their structure–composition–property relationships in detail. Moreover, as far as we know, no specific strategy has been reported for the design of antiperovskite alloys for photovoltaic applications.

In this work, we first investigated the electronic property, static dielectric constant, and exciton binding energy of six inorganic antiperovskites X<sub>3</sub>NA (X<sup>2+</sup> = Mg<sup>2+</sup>, Ca<sup>2+</sup>, Sr<sup>2+</sup>; A<sup>3-</sup> = P<sup>3-</sup>, As<sup>3-</sup>, Sb<sup>3-</sup>, Bi<sup>3-</sup>) to establish the structure–composition–property relationships. Based on the calculated results, a general linear relationship between the tolerance factor and various physical quantities is uncovered. Guided by the structure–composition–property relationship, we design a promising antiperovskite alloy Mg<sub>3</sub>NA<sub>0.5</sub>Bi<sub>0.5</sub> with a quasi-direct band gap of 1.402 eV. Finally, we make a comprehensive comparison between antiperovskite nitrides and conventional lead halide perovskites. Our studies will provide a strategy to design favorable antiperovskite alloys for novel device applications.

## ■ COMPUTATIONAL DETAILS

Our calculations are performed using the projector augmented wave (PAW) method implemented in the Vienna Ab initio Simulation Package (VASP) code.<sup>19,20</sup> The standard PAW pseudopotentials are adopted.<sup>21</sup> We use the Perdew, Burke, and Ernzerhof (PBE) form of the generalized gradient

approximation (GGA) exchange correlation functional<sup>22</sup> for crystal structure relaxation and the hybrid density functional (Heyd–Scuseria–Ernzerhof, HSE06)<sup>23</sup> for electronic properties of antiperovskites. The cutoff energy is set to 500 eV after convergence tests. We employ 3 × 3 × 3, 6 × 6 × 6, and 9 × 9 × 9  $\Gamma$ -centered Monkhorst–Pack *k*-point<sup>24</sup> grid for relaxations, self-consistent calculations, and density of states (DOSs) calculations, respectively. In our current calculations, the total energy is converged to less than 10<sup>-5</sup> eV. The maximum force is less than 0.02 eV/Å during the optimization, where both lattice constants and atomic positions are fully relaxed. The spin–orbit coupling (SOC) effects have been included for the electronic band structures of all of the studied antiperovskites X<sub>3</sub>BA. The phonon dispersion calculations are based on the PHONOPY code.<sup>25</sup> A 4 × 4 × 4 supercell with 2 × 2 × 2 *k* mesh is used to ensure the convergence.

For cubic antiperovskites, there are three nonzero elastic stiffness constants, C<sub>11</sub>, C<sub>12</sub>, and C<sub>44</sub>, and the stress–strain relationship is obtained from the Hooke’s law under the plane-stress condition  $E_i = C_{ij}\varepsilon_j$ . Here, we recalculate the strained sample using a higher cutoff energy and get the elastic constants C<sub>ij</sub> using the VASPKIT code.<sup>26</sup> Then, the Bulk modulus *B*, shear modulus *G*, Young’s modulus *Y*, and Poisson’s ratio  $\nu$  can be derived as<sup>27</sup>

$$B = \frac{C_{11} + 2C_{12}}{3}$$

$$G = \frac{C_{11} - C_{12} + 3C_{44}}{5}$$

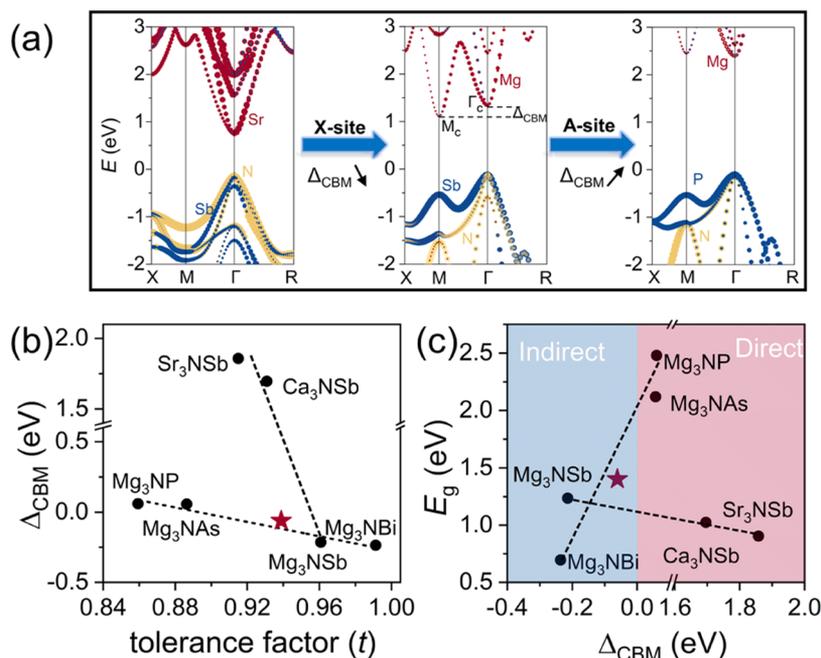
$$Y = \frac{9BG}{3B + G}$$

$$\nu = \frac{3B - 2G}{2(3B + G)}$$

The dielectric properties of the studied compounds are calculated by density functional perturbation theory (DFPT).<sup>28,29</sup> In DFPT, the dielectric constant tensor is defined as a linear response to the perturbative electric field and the ionic displacement is considered as a perturbation to

**Table 1.** Calculated Lattice Constant  $a$ , Goldschmidt's Tolerance Factor  $t$  (Table S1), Band Gap  $E_g$ , Static Dielectric Constant  $\epsilon_{\text{std}}$ , Exciton Binding Energy  $E_b$ , and Optical Transition Features between Band Edges at the  $\Gamma$  Point for  $X_3\text{NA}$  ( $X^{2+} = \text{Mg}^{2+}$ ,  $\text{Ca}^{2+}$ ,  $\text{Sr}^{2+}$ ;  $A^{3-} = \text{P}^{3-}$ ,  $\text{As}^{3-}$ ,  $\text{Sb}^{3-}$ ,  $\text{Bi}^{3-}$ ) Antiperovskites

compound	$a$ (Å)		$t$	$E_g$ (eV)	$\epsilon_{\text{std}}$	$E_b$ (meV)	optical transition
	calc.	expt.					
$\text{Mg}_3\text{NP}$	4.178		0.859	2.480	45.399	65	allowed
$\text{Mg}_3\text{NAs}$	4.236	4.217 <sup>35</sup>	0.886	2.119	36.130	12	allowed
$\text{Mg}_3\text{NSb}$	4.375	4.352 <sup>35</sup>	0.961	1.235	31.567	10	allowed
$\text{Mg}_3\text{NBi}$	4.437		0.991	0.696	38.919	4	allowed
$\text{Sr}_3\text{NSb}$	5.214	5.173 <sup>36</sup>	0.915	0.905	41.695	14	allowed
$\text{Ca}_3\text{NSb}$	4.873	4.854 <sup>37</sup>	0.931	1.025	32.752	12	allowed



**Figure 2.** (a) Schematic diagram of the band structure by changing the X- and A-site in antiperovskites  $X_3\text{NA}$  ( $X^{2+} = \text{Mg}^{2+}$ ,  $\text{Ca}^{2+}$ ,  $\text{Sr}^{2+}$ ;  $A^{3-} = \text{P}^{3-}$ ,  $\text{As}^{3-}$ ,  $\text{Sb}^{3-}$ ,  $\text{Bi}^{3-}$ ). Two lowest conduction band states  $M_c$  and  $\Gamma_c$  and the energy difference between the two states ( $\Delta_{\text{CBM}}$ ) are labeled. (b)  $\Delta_{\text{CBM}}$  as a function of tolerance factor, and (c) band gap varies as  $\Delta_{\text{CBM}}$  for cubic  $Pm\bar{3}m$  phase  $X_3\text{NA}$  antiperovskites. Red stars in (b) and (c) show the corresponding data for the  $\text{Mg}_3\text{NAs}_{0.5}\text{Bi}_{0.5}$  alloy.

**Table 2.** Percent Contributions from Each Atomic Orbital to the Two Conduction Band States  $M_c$  and  $\Gamma_c$  in  $X_3\text{NSb}$  ( $X^{2+} = \text{Sr}^{2+}$ ,  $\text{Ca}^{2+}$ ,  $\text{Mg}^{2+}$ )

compound	$\Delta_{\text{CBM}}$ (eV)	$M_c$			$\Gamma_c$		
		X (%)	Sb (%)	N (%)	X (%)	Sb (%)	N (%)
$\text{Sr}_3\text{NSb}$	1.858	50(s) + 13(p) + 21(d)	10(d)	6(s)	13(s) + 82(d)	5(d)	0
$\text{Ca}_3\text{NSb}$	1.696	44(s) + 7(p) + 30(d)	11(d)	8(s)	11(s) + 83(d)	6(d)	0
$\text{Mg}_3\text{NSb}$	-0.214	63(s) + 2(p)	10(d)	25(s)	70(s) + 13(d)	17(d)	0

the equilibrium system. Subsequently, the static dielectric constant ( $\epsilon_{\text{std}}$ ) consists of the electronic ( $\epsilon_{\text{ele}}$ ) and the ionic parts ( $\epsilon_{\text{ion}}$ ) of the system to the applied electric field.<sup>30</sup> The  $\epsilon_{\text{ele}}$  is the response of the electronic charge density to the perturbative electric field in the linear response regime. To increase the accuracy of dielectric constant calculations, the total energy is converged to less than  $10^{-6}$  eV, and the maximum force is less than 0.001 eV/Å.

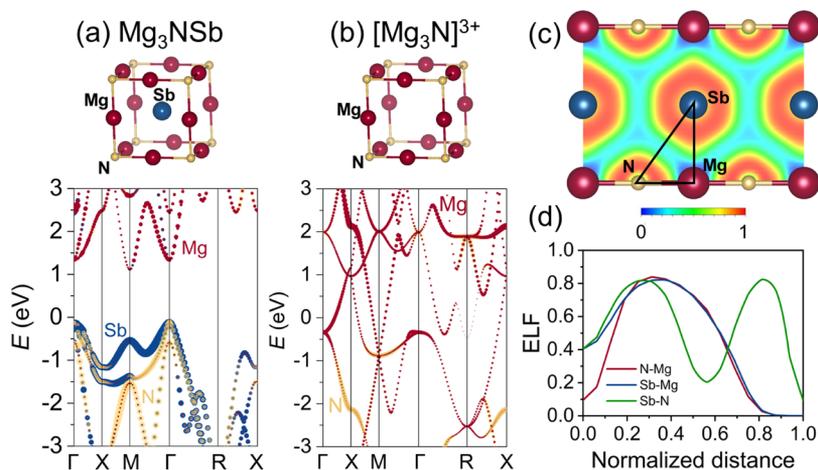
Based on the calculated  $\epsilon_{\text{ele}}$ , we estimate the exciton binding energy  $E_b$  using the Wannier model<sup>31</sup>

$$E_b = R_y \frac{\mu^*}{\epsilon^2}$$

where  $R_y = 13.56$  eV is the atomic Rydberg energy,  $\mu^*$  is the reduced exciton mass ( $1/\mu^* = 1/m_e + 1/m_h$ ), and  $\epsilon$  is the electronic dielectric constant.

## RESULTS AND DISCUSSION

**Effects of X- and A-Site Elements on Electronic Properties.** For conventional  $\text{ABX}_3$  perovskites, it is well known that the B-site cation and X-site anion have major impacts on the electronic structures, while the A-site cation does not contribute to the band edge because of its highly ionic nature.<sup>32,33</sup> However, the effect of composition on the electronic properties in antiperovskites has not been fully understood.<sup>34</sup> Here, we take six antiperovskite nitrides  $X_3\text{NA}$



**Figure 3.** Comparison of band structures of antiperovskites  $\text{Mg}_3\text{NSb}$  (a) and a frame model of  $[\text{Mg}_3\text{N}]^{3+}$  (b) with the same cubic lattice. The VBM (Fermi level) is shifted to zero for semiconductors (metal). (c) Electron localization function (ELF) in the (110) plane of cubic  $\text{Mg}_3\text{NSb}$ . (d) Line profile of ELF along three bonds, and the isosurface level is  $0.983 \text{ e}/\text{bohr}^3$ .

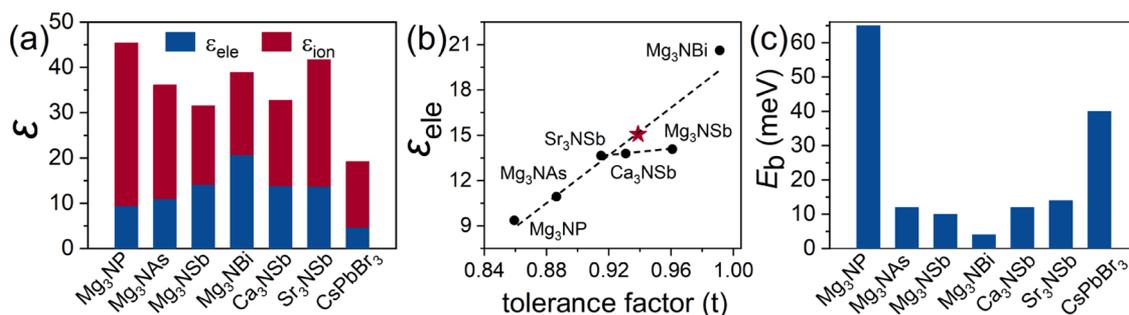
( $\text{X}^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ;  $\text{A}^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ) as examples to study their composition–property relationship. The stability, optimized structure parameters, and electronic properties are summarized in Table 1 and Supporting Information (see Figures S1–S6 and Table S2), which are consistent with previous reports. Here, we mainly focus on tuning the convergence of the lowest unoccupied state between M and  $\Gamma$  points to realize a direct–indirect/indirect–direct transition (see Figure 2a), which has been ignored in previous literatures. As shown in Figure 2a, we define quantitatively the energy difference between the two conduction band states  $\Gamma_c$  and  $M_c$  as  $\Delta_{\text{CBM}}$ . The positive value of  $\Delta_{\text{CBM}}$  means the direct band gap, and the negative value means the indirect band gap. When changing the X-site element from Sr to Ca to Mg, the value of  $\Delta_{\text{CBM}}$  decreases from 1.858 eV in  $\text{Sr}_3\text{NSb}$  to 1.696 eV in  $\text{Ca}_3\text{NSb}$ , to  $-0.214 \text{ eV}$  in  $\text{Mg}_3\text{NSb}$  (Table 2), meaning a direct–indirect transition. The change of  $\Delta_{\text{CBM}}$  also exists in  $\text{Mg}_3\text{NA}$  ( $\text{A}^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ) by substituting A-site elements. When replacing  $\text{Sb}^{3-}$  with a heavier element (i.e.,  $\text{Bi}^{3-}$ ),  $|\Delta_{\text{CBM}}|$  becomes larger. However, when  $\text{Sb}^{3-}$  is replaced with lighter elements (i.e.,  $\text{As}^{3-}$  and  $\text{P}^{3-}$ ), the value of  $\Delta_{\text{CBM}}$  changes from negative to positive, suggesting that the direct band gap is restored.

Taking  $\text{X}_3\text{NSb}$  ( $\text{X}^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ) as an example, we analyzed the reason behind the direct–indirect transition. The variation of the two states by the orbital contribution of the two states is shown in Table 2. For  $M_c$  states, the X-site contribution changes from higher energy Sr/Ca d orbitals in  $\text{X}_3\text{NSb}$  ( $\text{X}^{2+} = \text{Sr}^{2+}$  and  $\text{Ca}^{2+}$ ) to lower Mg 3s orbitals in  $\text{Mg}_3\text{NSb}$ . As a result, the  $M_c$  state is pulled down as the X-site element changes from Sr to Ca to Mg. While the higher  $\Gamma_c$  state is caused by the enhanced orbital overlap between  $\text{X}^{3-}$  and  $\text{Sb}^{3-}$  in  $\text{Mg}_3\text{NSb}$ , which is consistent with the partial DOS (PDOS) in Figure S5. This heightened orbital overlap generally correlates with the smaller X–Sb bond length. The reason for the indirect–direct transition induced by the change of the A-site composition is similar. Interestingly, this similar direct to indirect transition has also been observed in two-dimensional crystals from  $\text{Si}_3\text{O}$  to  $\text{C}_3\text{O}$ .<sup>38</sup>

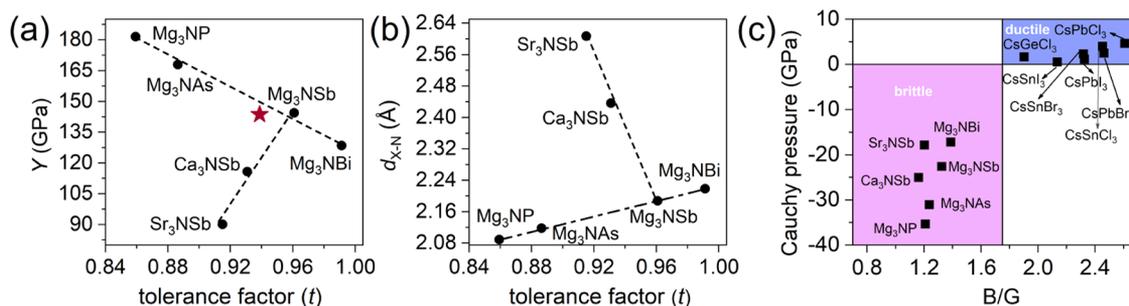
To rationalize the changing trend of the  $\Delta_{\text{CBM}}$  in  $\text{X}_3\text{NA}$  ( $\text{X}^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ;  $\text{A}^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ), we summarize the  $\Delta_{\text{CBM}}$  value as a function of the Goldschmidt's tolerance

factor ( $t$ ) in Figure 2b, because the  $t$  is an important geometric parameter to describe the structure stability of perovskite materials. Interestingly, regardless of the replacement of X-site or A-site, the  $\Delta_{\text{CBM}}$  shows a negative linear dependence on the  $t$ , as a consequence of lower s orbitals of Mg and Bi elements. Therefore, we can adjust the values of  $\Delta_{\text{CBM}}$  (e.g., from  $\Delta_{\text{CBM}} < 0 \text{ eV}$  to  $> 0 \text{ eV}$ ) to optimize the corresponding indirect/direct band gap feature through choosing a suitable  $t$ . We further establish the relationship between the band gap and the values of  $\Delta_{\text{CBM}}$  in antiperovskites  $\text{X}_3\text{NA}$  ( $\text{X}^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ;  $\text{A}^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ), as shown in Figure 2c. It is shown that the band gap ( $E_g$ ) decreases linearly with the  $\Delta_{\text{CBM}}$  when changing the X-site in group IIA but increases linearly with the  $\Delta_{\text{CBM}}$  when changing the A-site in group VA. The positive  $\Delta_{\text{CBM}}$  and optimal band gap ( $\sim 1.5 \text{ eV}$ ) are desired for single-junction solar cells. Therefore, we believe that the  $\Delta_{\text{CBM}}-E_g$  map diagram can serve as an important guidance for designing antiperovskite alloys in the future. Meanwhile, when changing X- or A-site elements, a linear relationship between the tolerance factor and band gap, effective mass, and valence band width is also observed in the studied nitride antiperovskites (see Figure S7). This is because the s/d levels of the X-site go down in energy from  $\text{Mg}^{2+}$  to  $\text{Sr}^{2+}$  (see Figure S8), pulling down the CBM and reducing the band gap. We note that a similar linear relationship has already been theoretically and experimentally reported in traditional inorganic halide perovskite alloys.<sup>39,40</sup>

From the change of  $\Delta_{\text{CBM}}$  in  $\text{Mg}_3\text{NA}$  ( $\text{A}^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ), it can be observed that the A-site does contribute to the electronic structure (see the PDOS in Figure S6), which is very different from the conventional halide perovskites. To further clarify the key role of the A-site in the electronic structure, we construct a hypothetical frame model  $[\text{Mg}_3\text{N}]^{3+}$  with  $\text{Sb}^{3-}$  ions removed in the same crystal lattice (see Figure 3) and calculate its band structure. Surprisingly, the system turns out to be a metal after removing all Sb atoms, which is quite different from the original semiconducting  $\text{Mg}_3\text{NSb}$ , further confirming the significant contributions of A-sites to the electronic structure in antiperovskites. On the contrary, in conventional halide perovskites, the band structure of  $[\text{PbI}_3]^-$  model without  $\text{Cs}^+$  ions is almost the same as that of the original  $\text{CsPbI}_3$  (see Figure S9), indicating that the A-site cations do not affect the



**Figure 4.** (a) Ionic ( $\epsilon_{\text{ion}}$ ) and electronic ( $\epsilon_{\text{ele}}$ ) contributions to the static dielectric constant ( $\epsilon_{\text{std}}$ ) for  $X_3\text{NA}$  ( $X^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ;  $A^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ). (b) Electronic dielectric constant ( $\epsilon$ ) as a function of the tolerance factor for  $X_3\text{NA}$ ; the red star shows the corresponding data for the  $\text{Mg}_3\text{NAS}_{0.5}\text{Bi}_{0.5}$  alloy. (c) Exciton binding energy ( $E_b$ ) of  $X_3\text{NA}$ . The corresponding parameters in  $\text{CsPbBr}_3$  are also listed for comparison.<sup>42</sup>



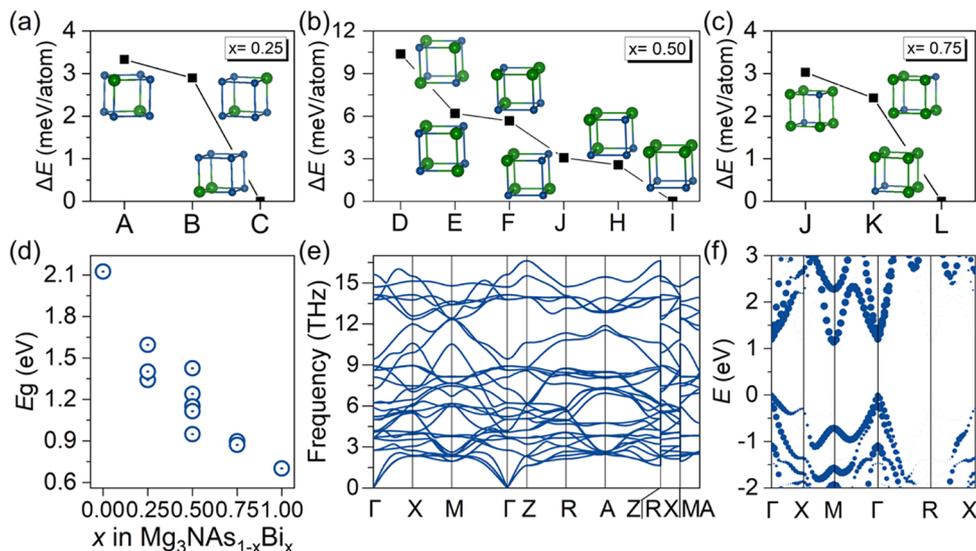
**Figure 5.** (a) Young's modulus  $Y$  (GPa) and (b) bond length of  $X\text{-N}$   $d_{X\text{-N}}$  (Å) as functions of the tolerance factor for antiperovskites  $X_3\text{NA}$  ( $X^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ;  $A^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ). The red star in (a) shows the corresponding data for the  $\text{Mg}_3\text{NAS}_{0.5}\text{Bi}_{0.5}$  alloy. (c) Correlation between Cauchy pressure and  $B/G$  for all of the studied six nitride antiperovskites  $X_3\text{NA}$  ( $X^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ;  $A^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ). The data for perovskites  $\text{CsBX}_3$  ( $B^{2+} = \text{Pb}^{2+}, \text{Sn}^{2+}, \text{Ge}^{2+}$ ;  $X^- = \text{Cl}^-, \text{Br}^-, \text{I}^-$ ) have been collected for comparison.

electronic states around the band edge. Further, from the electron localization function (ELF), it can be seen that the ELF of the  $A\text{-X/B-X}$  bond in the region between the nuclei of the atoms in  $\text{Mg}_3\text{NSb}$  (i.e.,  $\text{Sb-Ca}$  or  $\text{Sb-N}$ ,  $\text{ELF} > 0.2$ ) is significantly larger than that of  $\text{CsPbI}_3$  (i.e.,  $\text{Cs-I}$  or  $\text{Cs-Pb}$ ,  $\text{ELF} < 0.1$  in Figure S10). This implies that the  $A$ -site forms a strong bonding with other sites in the antiperovskites and, thus, has a non-negligible contribution to electronic properties accordingly.

**Static Dielectric Constants.** In addition to band gap, static dielectric constant is also an important descriptor for evaluating the performance of a solar cell absorber. For example, a large static dielectric constant provides strong screening, thereby suppressing carrier scattering, trapping, and recombination, eventually improving transport properties.<sup>41</sup> Although the electronic properties of antiperovskites have been widely reported, the dielectric properties of most of them have not yet been reported. We next focus on the static dielectric constants ( $\epsilon$ ) of  $X_3\text{NA}$  ( $X^{2+} = \text{Mg}^{2+}, \text{Ca}^{2+}, \text{Sr}^{2+}$ ;  $A^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ) and summarize the results in Table S3 and Figure 4. The  $\epsilon$  values range from 31.57 in  $\text{Mg}_3\text{NSb}$  to 45.40 in  $\text{Mg}_3\text{NP}$ , and the data of  $\text{Mg}_3\text{NSb}$  are in line with a previous report.<sup>16</sup> All of the  $\epsilon$  values are substantially enhanced compared with the calculated values in halide perovskites (i.e.,  $\text{CsPbBr}_3$ ,  $\sim 20.00$ ) (see Figure 4a and Table S3),<sup>42</sup> which can be attributed to the larger ionic ( $\epsilon_{\text{ion}}$ ) and electronic ( $\epsilon_{\text{ele}}$ ) contributions. Here, we mainly focus on the  $\epsilon_{\text{ele}}$  that is related to the electronic properties directly. It is shown that the  $\epsilon_{\text{ele}}$  in antiperovskites  $X_3\text{NA}$  is at least twice as large as those in halide perovskites. This may be due to the significant contribution of the  $A$ -site anion to the electronic properties of antiperovskites, which has been discussed before. Further, we find a positive

linear relationship between  $\epsilon_{\text{ele}}$  and  $t$  in  $\text{Mg}_3\text{NA}$  ( $A^{3-} = \text{P}^{3-}, \text{As}^{3-}, \text{Sb}^{3-}, \text{Bi}^{3-}$ ), as shown in Figure 4b. With the increase of  $t$ , the  $\epsilon_{\text{ele}}$  increases from 9.37 at  $t = 0.86$  in  $\text{Mg}_3\text{NP}$  to 20.63 at  $t = 0.99$  in  $\text{Mg}_3\text{NBI}$ . This is related to the linear relationship between band gap and tolerance factor, as shown in Figure S7, because a larger band gap generally leads to a smaller  $\epsilon_{\text{ele}}$ .<sup>43</sup> On the other hand, we find that the  $\epsilon_{\text{ele}}$  is not sensitive to the  $X$ -site element because of their similar band gap values. Based on the calculated  $\epsilon_{\text{ele}}$ , we estimate the exciton binding energies  $E_b$  of the six studied antiperovskites  $X_3\text{NA}$  (see Table 1 and Figure 4c) using the Wannier model, which is sufficient to reflect the correct trends of the  $E_b$  values (see Figure S11).  $\text{Mg}_3\text{NP}$  has the largest  $E_b$  of 65 meV, as a result of its largest effective mass and smallest electronic dielectric constant. The estimated  $E_b$  for other five  $X_3\text{NA}$  ranges from 4 to 14 meV (see Table 1 and Figure 4c), smaller than the thermal energy ( $k_B T \sim 26$  meV) at room temperature. These  $E_b$  values are only 30% of those in star lead halide perovskites (i.e.,  $\text{CsPbBr}_3$ ). The small exciton binding energy in these antiperovskites will facilitate rapid electron-induced carrier dissociation, which is considerably important for an ideal solar cell absorber.

**Mechanical Properties.** For practical fabrication and device applications, the mechanical properties of materials are very important and they can be useful to predict their aging behaviors. We, thus, summarize the calculated elastic constants of all of the studied antiperovskites  $X_3\text{NA}$  in Table S4. The dependence of the derived Young's modulus ( $Y$ ) on the tolerance factor ( $t$ ) is shown in Figure 5a. A similar linear relationship is clearly identified between  $Y$  and  $t$ . For instance, for the Mg-based nitride compounds, the  $Y$  value declines monotonically with increasing  $t$  when changing the  $A$ -site in group VA. The  $Y$  value decreases from 181.48 GPa at  $t = 0.86$



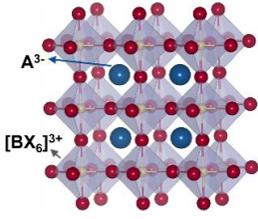
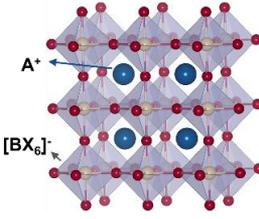
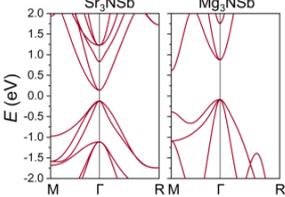
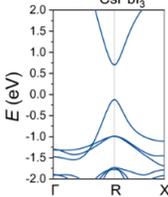
**Figure 6.** Energies of  $\text{Mg}_3\text{NAS}_{1-x}\text{Bi}_x$  with different types of As (blue) and Bi (green) motif arrangements for (a)  $x = 0.25$ , (b)  $x = 0.50$ , and (c)  $x = 0.75$ . The energies of the lowest energy configurations at each concentration are set to zero. (d) Calculated band gaps of  $\text{Mg}_3\text{NAS}_{1-x}\text{Bi}_x$  as a function of  $x$ , the amount of Bi. (e) Phonon dispersions and (f) unfolded band structures of the  $\text{Mg}_3\text{NAS}_{0.5}\text{Bi}_{0.5}$  alloy in the most stable configuration.

in  $\text{Mg}_3\text{NP}$  to 128.49 GPa at  $t = 0.99$  in  $\text{Mg}_3\text{NBI}$ . However, for the  $\text{X}_3\text{NSb}$ , the  $Y$  value increases monotonically with increasing  $t$ . When replacing the  $\text{X}^-$  site in group IIA, it ranges from 90.10 GPa at  $t = 0.915$  in  $\text{Sr}_3\text{NSb}$  to 144.46 GPa at  $t = 0.961$  in  $\text{Mg}_3\text{NSb}$ . This linear variation of Young's modulus can be related to the  $\text{N-X}$  bond length in  $\text{X}_3\text{NA}$  as shown in Figure 5b, where a smaller bond length generally requires a large strain to break out, leading to a large  $Y$ . Next, we elaborate the ductile or brittle nature of these antiperovskites using Cauchy pressure, which is defined as the difference between the two particular elastic constants  $C_{12} - C_{44}$ , serving as an indication of ductility. The material is expected to be brittle (ductile) for negative (positive) Cauchy pressure. Here, the Cauchy pressure for antiperovskite nitrides  $\text{X}_3\text{NA}$  is negative, clearly showing the brittle nature of these compounds. The ratio between Bulk modulus  $B$  and shear modulus  $G$  is another index associated with the ductile nature of materials. A low (high)  $B/G$  ratio suggests the brittle (ductile) feature of compounds, and the critical value is 1.75.<sup>44</sup> This ratio ranges from 1.1 to 1.5 for the six  $\text{X}_3\text{NA}$ , further confirming the brittle nature of antiperovskites. The brittle nature of antiperovskites  $\text{X}_3\text{NA}$  is very different from the ductile feature of common halide perovskites  $\text{CsBX}_3$  ( $\text{B}^{2+} = \text{Pb}^{2+}, \text{Sn}^{2+}, \text{Ge}^{2+}$ ;  $\text{X}^- = \text{Cl}^-, \text{Br}^-, \text{I}^-$ ), which is clearly indicated in the map diagram in Figure 5c. Therefore, the cubic antiperovskite nitrides  $\text{X}_3\text{NA}$  may possess a low probability of application in flexible devices in the future.

**Designing Promising Antiperovskite through A-Site Alloying.** The alloy strategy has been proven to be an effective way to modify the properties of perovskites and overcome some disadvantages of single compounds,<sup>45</sup> and the above established structure–composition–property relationship is very instructive for the design of antiperovskite alloys. Here, we investigate the electronic properties of  $\text{Mg}_3\text{NAS}_{1-x}\text{Bi}_x$  alloys with  $x = 0.25, 0.50,$  and  $0.75$  for the first time, which may possess an ideal band gap ( $\sim 1.5$  eV) for single-junction solar cells based on the linear relationship, as shown in Figure 2c. The cubic random alloys are modeled using  $2 \times 2 \times 2$  supercells containing 40 atoms. The Bi atoms can substitute As

atoms at different positions, thus forming several different configurations with certain  $x$ . For example, there are 3, 6, and 3 different configurations for  $x = 0.25, 0.50,$  and  $0.75$  (see Figure 6a–c), respectively. The total energy difference between different configuration is smaller than 11 meV/atom, suggesting the possibility of forming disordered solid solutions.<sup>46</sup> It is shown that although the band gap differs between different configurations for a certain  $x$ , the band gap of  $\text{Mg}_3\text{NAS}_{1-x}\text{Bi}_x$  alloys generally decreases with increasing  $x$ , reduced from 2.102 eV in  $\text{Mg}_3\text{NAS}$  to 0.612 eV in  $\text{Mg}_3\text{NBI}$  (see Figure 6d). This decreasing trend is associated with the higher Bi p orbitals than As p orbitals.

Finally, to validate the practicability of the structure–composition–property relationship, we focus on  $\text{Mg}_3\text{NAS}_{0.5}\text{Bi}_{0.5}$  in the most stable configuration I, which has been further verified by a statistical approach (see Figure S12). The configuration I shows no imaginary phonon frequencies in Figure 6e, indicating kinetic stability of the structure I. The tolerance factor of this alloy is 0.94. Figure 6f shows the band structure with the bands unfolded back from supercell to its primitive unit cell, which can be comparable with the bands of primitive cell, and quantitative analysis of the alloy is enabled.<sup>47</sup> It clearly shows a quasi-direct band gap of 1.402 eV. This data has been added as a red star sign in the linear relationship in Figure S7, located at the fitted line, confirming the linear rule in the band gap again. This is another confirmation that we can use this rule to design antiperovskites with an optimal band gap ( $\sim 1.5$  eV) for solar cells. Furthermore, the  $\Delta_{\text{CBM}}$  is  $-0.061$  eV, agreeing well with the predicted data of  $-0.059$  eV using the linear relationship in Figure 2b. For dielectric constants of  $\text{Mg}_3\text{NAS}_{0.5}\text{Bi}_{0.5}$ , they are isotropic in  $x$  and  $y$  directions but a little anisotropic in  $z$  direction. The  $\epsilon_{\text{ele}}$  ( $\epsilon_{\text{ion}}$ ) is around 15.06 (23.00), as shown in Table S3. Finally, the Young's modulus for  $\text{Mg}_3\text{NAS}_{0.5}\text{Bi}_{0.5}$  is 143.61 GPa. Both the  $\epsilon_{\text{ele}}$  and  $Y$  are on the fitting lines in Figures 4b and 5a. Thus, the physical quantities of the  $\text{Mg}_3\text{NAS}_{0.5}\text{Bi}_{0.5}$  alloy confirm the uncovered structure–composition–property relationship. We can use this universal

	3D antiperovskite X <sub>3</sub> BA	3D halide perovskite ABX <sub>3</sub>
Crystal structure		
Material example	A <sup>3-</sup> = P <sup>3-</sup> , As <sup>3-</sup> , Sb <sup>3-</sup> , Bi <sup>3-</sup> B <sup>3-</sup> = N <sup>3-</sup> X <sup>2+</sup> = Mg <sup>2+</sup> , Ca <sup>2+</sup> , Sr <sup>2+</sup>	A <sup>+</sup> = MA <sup>+</sup> , FA <sup>+</sup> , Cs <sup>+</sup> B <sup>2+</sup> = Pb <sup>2+</sup> X <sup>-</sup> = Cl <sup>-</sup> , Br <sup>-</sup> , I <sup>-</sup>
Synthesis	Reaction tube (~800 °C)	In solution (~100 °C)
Processing methods	sintering, sputtering	CVD, sputtering
Sample types	Powder, thin film	Single crystal, thin film
Typical band structures		
Band edge's main contribution	CBM X <i>d/s</i> VBM A <i>p</i> + B <i>p</i>	CBM B <i>p</i> VBM B <i>s</i> + X <i>p</i>
<i>E<sub>g</sub></i> (eV)	0.70-2.48	1.40-3.21 <sup>48, 49</sup>
<i>m<sub>h</sub></i> ( <i>m</i> <sub>0</sub> )	0.26-0.98	0.14-0.31 <sup>48, 50</sup>
<i>E<sub>b</sub></i> (meV)	4-65	14-68 <sup>51</sup>
Optical transition	Allowed	Allowed
<i>Y</i> (GPa)	90.10-181.48	17.79-21.92 <sup>52</sup>
Brittleness/Ductility	Brittleness	Ductility

**Figure 7.** Comprehensive comparison of 3D antiperovskite X<sub>3</sub>BA and the typical 3D halide perovskite ABX<sub>3</sub> for the band structure, electronic band gap *E<sub>g</sub>*, exciton binding energy *E<sub>b</sub>*, carrier effective mass *m<sub>h</sub>*, optical transition character, Young's modulus *Y*, softness, synthesis method, and sample type.<sup>48–52</sup>

relationship to design suitable antiperovskite alloys for different optoelectronic applications.

**Comparison with Lead Halide Perovskites.** Toward future practical applications, let us now make a comprehensive comparison between antiperovskite nitrides X<sub>3</sub>NA (X<sup>2+</sup> = Mg<sup>2+</sup>, Ca<sup>2+</sup>, Sr<sup>2+</sup>; A<sup>3-</sup> = P<sup>3-</sup>, As<sup>3-</sup>, Sb<sup>3-</sup>, Bi<sup>3-</sup>) and the widely studied three-dimensional (3D) lead halide perovskites APbX<sub>3</sub> (A<sup>+</sup> = MA<sup>+</sup>, FA<sup>+</sup>, Cs<sup>+</sup>; X<sup>-</sup> = Cl<sup>-</sup>, Br<sup>-</sup>, I<sup>-</sup>) (see Figure 7). Overall, X<sub>3</sub>NA possess very similar electronic and optical properties as those of APbX<sub>3</sub>, such as suitable band gaps (~1.5 eV), small carrier effect masses (0.26–0.98 *m*<sub>0</sub>), small exciton binding energies (4–65 meV), and allowed optical transitions at band edges.

The excellent photoelectric properties can be attributed to their high symmetric crystal lattice, antibonding states of the VBM (see Figure S13), and high orbital connectivity near the band edges. Compared with APbX<sub>3</sub>, the antiperovskites X<sub>3</sub>NA are obtained by electronically inverting the formula, resulting in totally different band-edge characteristics. More interestingly, the effect of the A-site element on the electronic

structure has a non-negligible effect in the antiperovskites, which is quite different from the halide perovskites. In addition, according to the predicted mechanical properties, antiperovskites are difficult to be applied in flexible devices in the future. On the other hand, compared with the simple low-temperature solution processing of halide perovskites, most of the antiperovskites synthesized in current experiments are powder samples, which are prepared at high temperatures, such as Mg<sub>3</sub>NAs and Mg<sub>3</sub>NSb.<sup>16,35</sup> Therefore, although the successful synthesis of the antiperovskites films in experiment has been reported, the preparation of high-quality films is still challenging, which is a key step for future device applications.

## CONCLUSIONS

In summary, we have studied the previously overlooked conduction band convergence, dielectric constant, and exciton binding energy of cubic *Pm3m* phase antiperovskites X<sub>3</sub>NA (X<sup>2+</sup> = Mg<sup>2+</sup>, Ca<sup>2+</sup>, Sr<sup>2+</sup>; A<sup>3-</sup> = P<sup>3-</sup>, As<sup>3-</sup>, Sb<sup>3-</sup>, Bi<sup>3-</sup>) using first-principles calculations. Similar to lead halide perovskites, the X<sub>3</sub>NA has suitable band gaps (~1.5 eV), small carrier

effective masses (0.26–0.98  $m_0$ ), small exciton binding energies (4–65 meV), and allowed optical transitions at band edges. On the other hand,  $X_3NA$  exhibits totally different band-edge characteristics compared to the perovskites with electronically inverting the formula. Both the X-site and A-site can effectively tune the conduction band convergence, leading to the transition between indirect and direct band gap features. Importantly, a universal relationship between the tolerance factor and physical quantities, including band gap,  $\Delta_{CBM}$ , electronic dielectric constant, and Young's modulus, is uncovered. The linear relationship originates from the atomic orbital energies of X- and A-site elements. Based on the established structure–composition–property relationship in six antiperovskite nitrides  $X_3NA$ , we design the alloy  $Mg_3NAs_{0.5}Bi_{0.5}$  with an optimal band gap of 1.402 eV as a solar cell absorber. Finally, we make a comprehensive comparison between the nitride-based antiperovskites and lead halide perovskites. Our work will provide an effective strategy for designing promising antiperovskite alloys for novel device applications.

## ■ ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsami.1c10137>.

Stability, crystal and electronic structures, mechanical properties, band structure unfolding in alloy calculations, COHP analysis, structure–composition–property relationships in other antiperovskite and perovskites, and the effects of B-site elements on antiperovskites (PDF)

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## Author Contributions

G.T. conceived the idea and designed the project. H.Z. carried out most of the calculations and analyzed all data under S.Y.'s supervision. C.F. calculated the dielectric and mechanical properties. H.Z. and G.T. wrote the manuscript. S.Y. and G.T. revised the manuscript. All authors contributed to the discussion and revision of the paper.

## Notes

The authors declare no competing financial interest.

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